

ELEC0041

Modelling and Design of Electromagnetic Systems

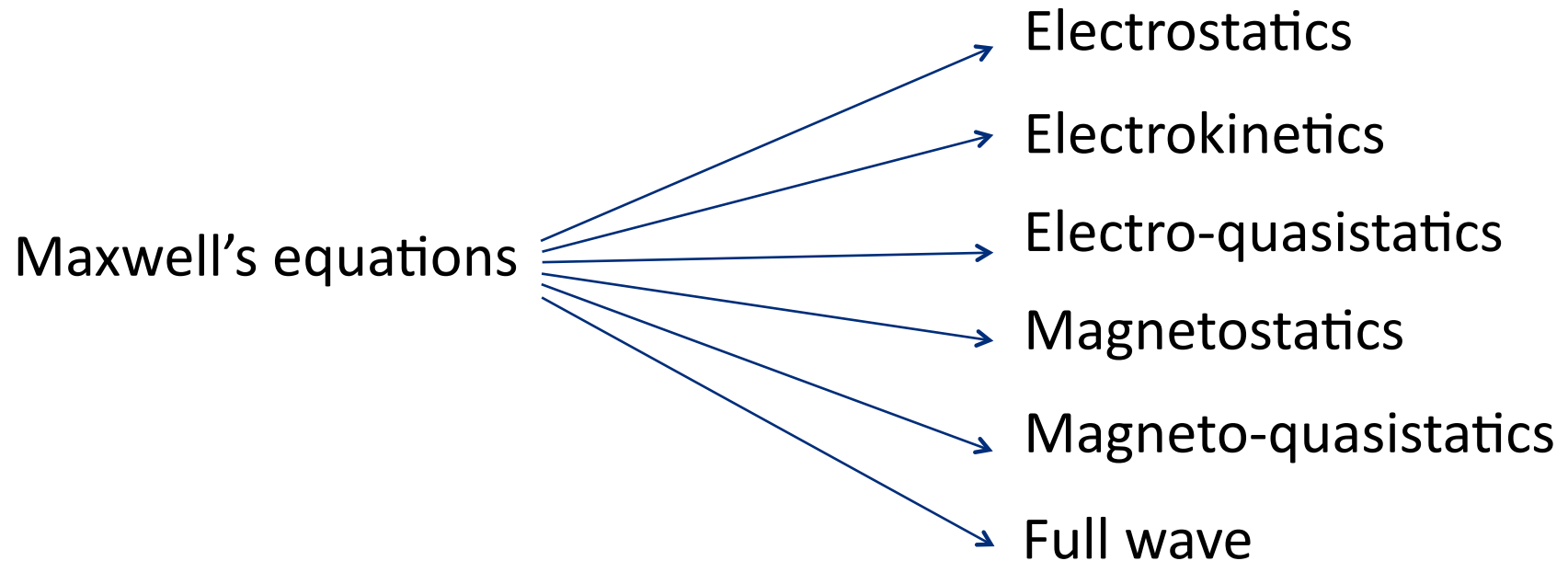
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Electromagnetic Models

Electromagnetic Models



Electromagnetic Models



- Electrostatics

- Distribution of electric field due to static charges and/or levels of electric potential



- Electrokinetics

- Distribution of static electric current in conductors



- Electro-quasistatics (or "electrodynamics")

- Distribution of electric field and electric current in materials (insulating and conducting)



- Magnetostatics

- Distribution of static magnetic field due to magnets and continuous currents



- Magneto-quasistatics (or "magnetodynamics")

- Distribution of magnetic field and eddy current due to moving magnets and time variable currents



- Full wave

- Propagation of electromagnetic fields

Maxwell's Equations

$$\mathbf{curl} \mathbf{h} = \mathbf{j} + \partial_t \mathbf{d}$$

Maxwell-Ampère's equation

$$\mathbf{curl} \mathbf{e} = -\partial_t \mathbf{b}$$

Faraday's equation

$$\left. \begin{aligned} \mathbf{div} \mathbf{b} &= 0 \\ \mathbf{div} \mathbf{d} &= \rho_q \end{aligned} \right\}$$

Conservation equations

\mathbf{h} magnetic field (A/m)

\mathbf{e} electric field (V/m)

\mathbf{b} magnetic flux density (T)

\mathbf{d} electric displacement (C/m²)

\mathbf{j} current density (A/m²)

ρ_q charge density (C/m³)

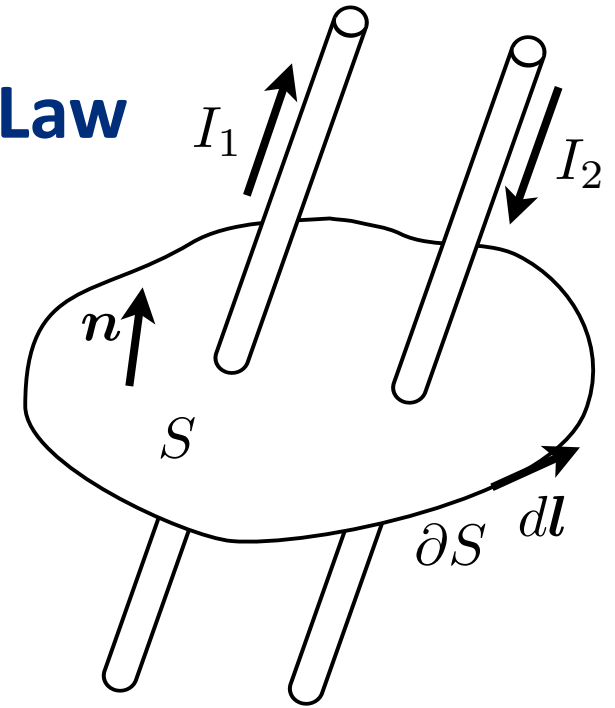
Integral form: Ampère's Law

$$\text{curl } \mathbf{h} = \mathbf{j}$$

$$\Rightarrow \oint_{\partial S} \mathbf{h} \cdot d\mathbf{l} = I$$

Magnetomotive force
(m.m.f.)

Circulation of magnetic field along closed curve equals algebraic sum of currents crossing the underlying surface



$$\oint_{\partial S} \mathbf{h} \cdot d\mathbf{l} = I_1 - I_2$$

Conservation of current: $\text{div } \mathbf{j} = 0$

$$\Rightarrow \oint_{\partial V} \mathbf{j} \cdot \mathbf{n} ds = 0$$

Sum of currents arriving at a node is zero

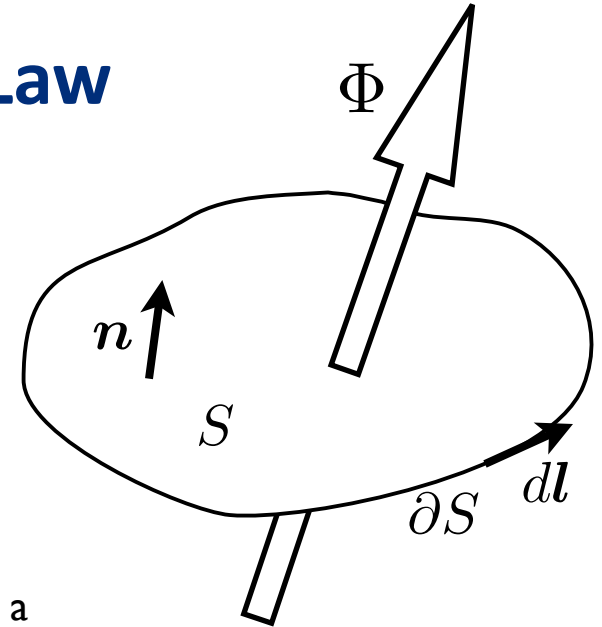
Integral form: Faraday's Law

$$\text{curl } \mathbf{e} = -\partial_t \mathbf{b}$$

$$\Rightarrow \oint_{\partial S} \mathbf{e} \cdot d\mathbf{l} = -\partial_t \Phi$$

Electromotive force
(e.m.f.)

Any variation of magnetic flux through a circuit gives rise to an electromotive force



For a circuit moving at speed v :

$$\text{f.e.m.} = \oint_{\partial S(t)} \mathbf{f}/q \cdot d\mathbf{l} = \oint_{\partial S(t)} (\mathbf{e} + \mathbf{v} \times \mathbf{b}) \cdot d\mathbf{l} = -d_t \int_{S(t)} \mathbf{b} \cdot \mathbf{n} ds$$

Conservation of magnetic flux density: $\text{div } \mathbf{b} = 0$

$$\Rightarrow \oint_{\partial V} \mathbf{b} \cdot \mathbf{n} ds = 0$$

Magnetic flux lines are closed

Lorentz Force

Interaction of electromagnetic fields with a point charge moving at speed v

$$\mathbf{f} = q(\mathbf{e} + \mathbf{v} \times \mathbf{b})$$

For a conductor (electrically neutral, only negative charges moving):

$$\mathbf{f} = \mathbf{j} \times \mathbf{b} = \mathbf{curl} \mathbf{h} \times \mathbf{b} \quad \text{Laplace Force}$$

Electromagnetic Power

Poynting vector: $\mathbf{s} = \mathbf{e} \times \mathbf{h}$

Power exchanged with a volume (interior normal):

$$P = \oint_{\partial V} \mathbf{s} \cdot \mathbf{n} \, ds = - \int_V \operatorname{div} \mathbf{s} \, dv = \int_V p \, dv$$

Power density:

$$p = -\operatorname{div} \mathbf{e} \times \mathbf{h} = -\mathbf{h} \cdot \operatorname{rot} \mathbf{e} + \mathbf{e} \cdot \operatorname{rot} \mathbf{h}$$

$$\Rightarrow p = \mathbf{h} \cdot \partial_t \mathbf{b} + \mathbf{e} \cdot \mathbf{j} + \mathbf{e} \cdot \partial_t \mathbf{d}$$

“magnetic”

“conduction”

“electric”

Material Constitutive Relations

$$\mathbf{b} = \mu \mathbf{h} \quad \text{Magnetic relation}$$

$$\mathbf{d} = \varepsilon \mathbf{e} \quad \text{Dielectric relation}$$

$$\mathbf{j} = \sigma \mathbf{e} \quad \text{Ohm's law}$$

Characteristics of materials:

μ magnetic permeability (H/m)

ε dielectric permittivity (F/m)

σ electric conductivity (S/m)



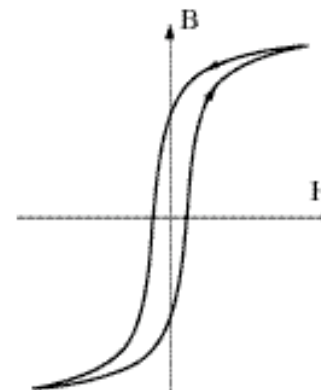
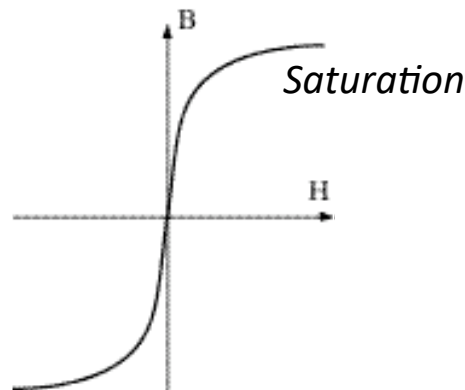
constants (linear materials),
 functions of electromagnetic fields (nonlinear
 materials), tensorial (anisotropic materials),
 functions of other physical fields
 (temperature, ...)

Magnetic Relation

$$\mathbf{b} = \mu \mathbf{h} \quad \mu = \mu_r \mu_0 \begin{cases} \mu_r & \text{Relative magnetic permeability} \\ \mu_0 & \text{Vacuum permeability (} 4\pi 10^{-7} \text{H/m)} \end{cases}$$

- Diamagnetic and paramagnetic materials $\mu_r \approx 1$
Linear materials (silver, copper, aluminum)
- Ferromagnetic materials $\mu_r \gg 1$, $\mu_r = \mu_r(h)$
Nonlinear materials (steel, iron)
Ferromagnetic-paramagnetic transition for $T > T_{\text{Curie}}$ (T_{Curie} of iron : 1043 K)

b(h) curves



Hysteresis (non bijective law)

Energy dissipation = cycle area

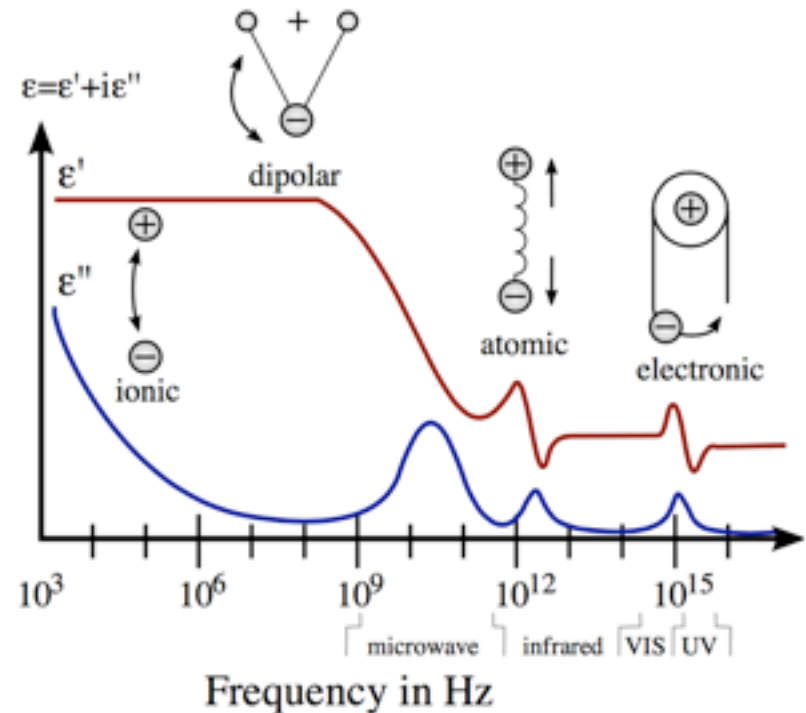
Steinmetz formula: $p_{\text{hyst}} = \omega k_h b_{\text{max}}^\nu$

Dielectric Relation

$$\mathbf{d} = \epsilon \mathbf{e} \quad \epsilon = \epsilon_r \epsilon_0 \quad \left\{ \begin{array}{l} \epsilon_r \text{ Relative dielectric permittivity} \\ \epsilon_0 = \frac{1}{\mu_0 c_0^2} \text{ Vacuum permittivity} \\ \quad \quad \quad (8.854187817620... \times 10^{-12} \text{ F/m}) \end{array} \right.$$

ϵ_r at room temperature for $f < 1\text{kHz}$

Air	1.0006
Teflon	2.1
Polyethylene	2.25
Paper	3.85
Glass	3.7 - 10
Concrete	4.5
Water	80



Ohm's Law

$$\mathbf{j} = \sigma \mathbf{e} \quad \left(\text{Resistivity } \rho = \frac{1}{\sigma} \right)$$

Simple models for temperature dependency

- Metals : $\rho = \rho_0(1 + \alpha(T - T_0))$

	$\rho_0 (T_0 = 20^\circ C)$ (Ωm)	α ($^\circ C^{-1}$)
Aluminum	$2.7 \cdot 10^{-8}$	$4 \cdot 10^{-3}$
Copper	$1.7 \cdot 10^{-8}$	$3.9 \cdot 10^{-3}$
Iron	$9.6 \cdot 10^{-8}$	$6.5 \cdot 10^{-3}$

- Glass : $\ln \rho = A + \frac{B}{T}$

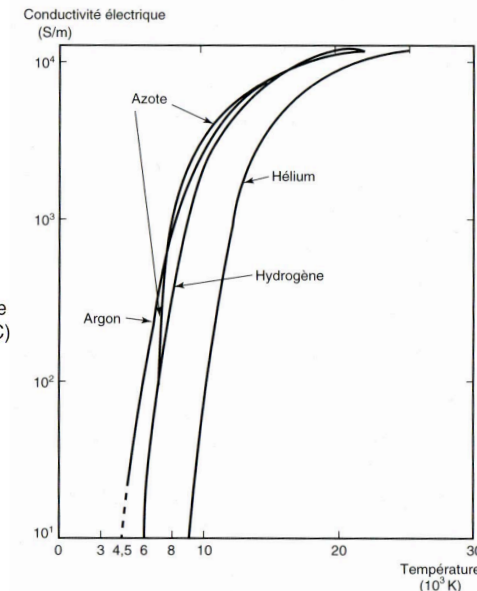
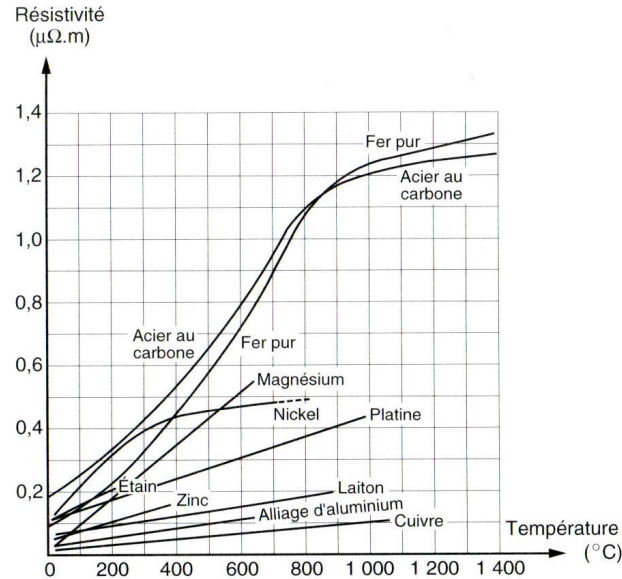
Common glass: $\ln \rho = -4.6 + \frac{7678}{T}$

- Ionic solutions : $\sigma = \sigma_0 + \alpha(T - T_0)$

Tap water: $\sigma_0 = 0.055 \Omega^{-1} m^{-1}$

$\alpha = 1.65 \cdot 10^{-3} \text{ } ^\circ C^{-1} \Omega^{-1} m^{-1}$

$T_0 = 20^\circ C$



Model Choice

Maxwell's equations & constitutive relations in frequency domain, without sources:

$$\Delta \mathbf{e} - i\omega\sigma\mu\mathbf{e} + \omega^2\varepsilon\mu\mathbf{e} = 0$$

Using characteristic lengths

- domain size L

- skin depth $\delta = \sqrt{\frac{2}{\omega\sigma\mu}}$

- wavelength $\lambda = \frac{2\pi}{k}$, with

$$\left\{ \begin{array}{l} \text{wave number } k = \frac{\omega}{c} \\ \text{speed of light } c = \frac{1}{\sqrt{\varepsilon\mu}} \end{array} \right.$$

allows to write in non-dimensional form:

$$\left(\frac{3}{L^2} - \frac{2i}{\delta^2} + \frac{4\pi}{\lambda^2} \right) \mathbf{e} = 0$$

Model Choice

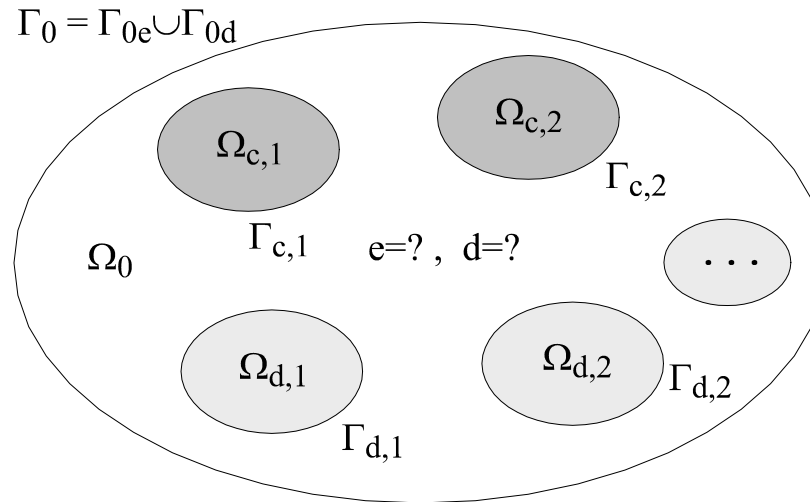
$$\left(\frac{3}{L^2} - \frac{2i}{\delta^2} + \frac{4\pi}{\lambda^2} \right) \mathbf{e} = 0 \quad \text{Non-dimensional numbers}$$

$$\left\{ \begin{array}{l} g_1 = \left(\frac{\lambda}{L} \right)^2 \\ g_2 = \left(\frac{\delta}{L} \right)^2 \\ g_3 = \left(\frac{\lambda}{\delta} \right)^2 \end{array} \right.$$

- $g_1 \gg 1$ uncoupled electric or magnetic problems
 - $g_2 \gg 1$ magnetostatics
 - $g_2 \lesssim 1$ magnetodynamics
 - $g_3 \gg 1$ electrokinetics
 - $g_3 \approx 1$ electrodynamics
 - $g_3 \ll 1$ electrostatics
- $g_1 \lesssim 1$ full wave ($g_1 \rightarrow 0$ high-frequency asymptotics)

Electrostatics

$$\begin{aligned} \text{curl } \mathbf{e} &= 0 \\ \text{div } \mathbf{d} &= \rho_q \\ \mathbf{d} &= \varepsilon \mathbf{e} \end{aligned}$$

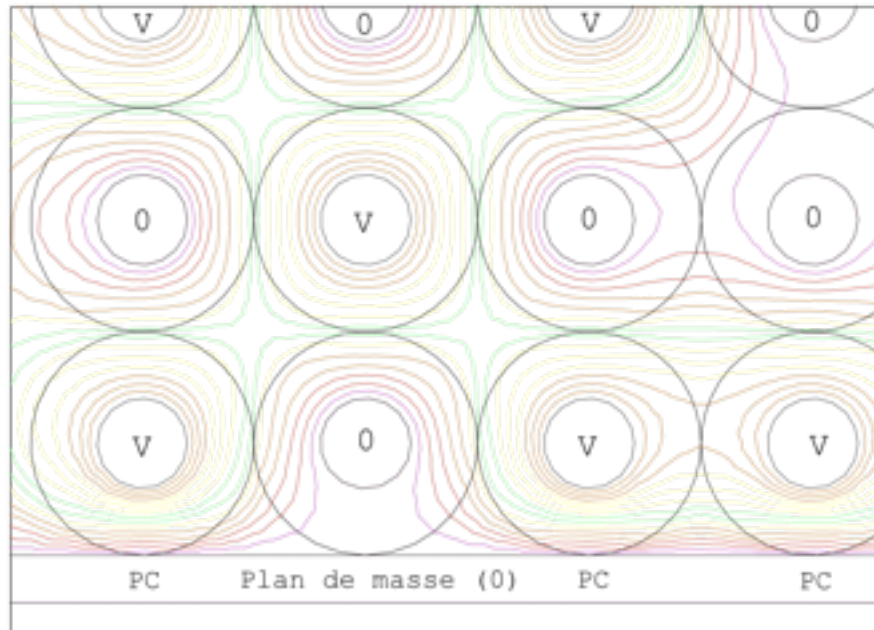


Example: electric scalar potential formulation

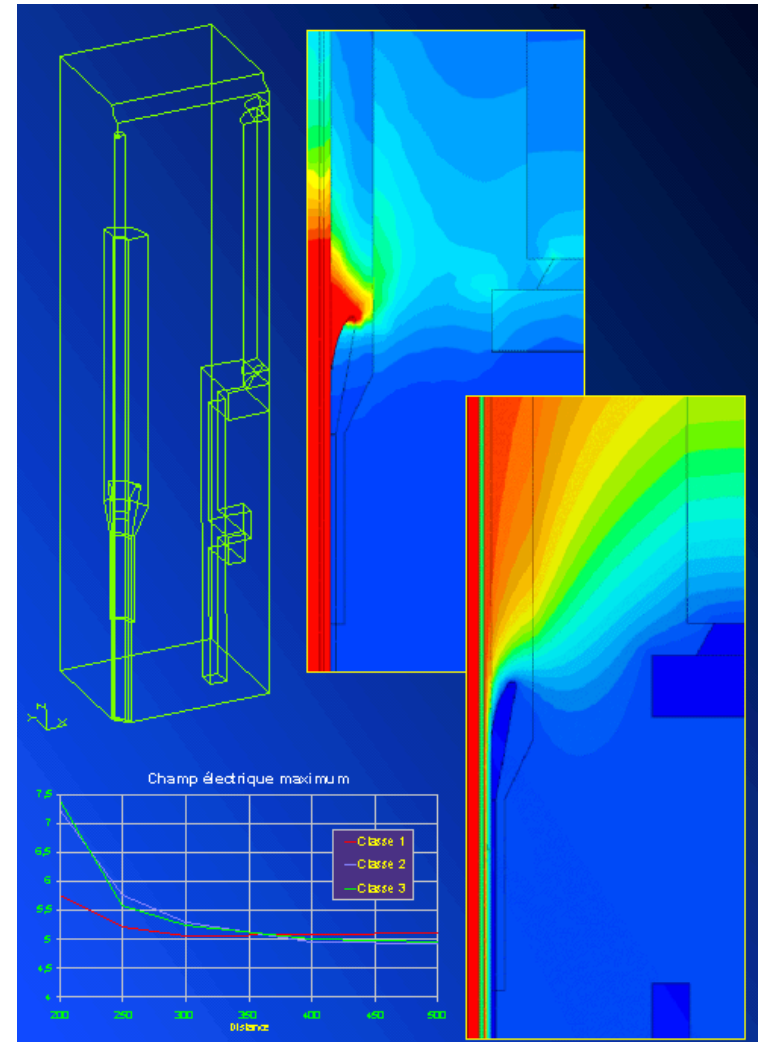
$$\text{div } \varepsilon \mathbf{grad } v = -\rho_q \quad \text{with} \quad \mathbf{e} = -\mathbf{grad } v$$

- Formulation for
 - the exterior region Ω_0
 - the dielectric regions $\Omega_{d,i}$
- In each conducting region $\Omega_{c,i}$, $v = v^i \Rightarrow v|_{\Gamma_{c,i}} = v^i$

Electrostatics

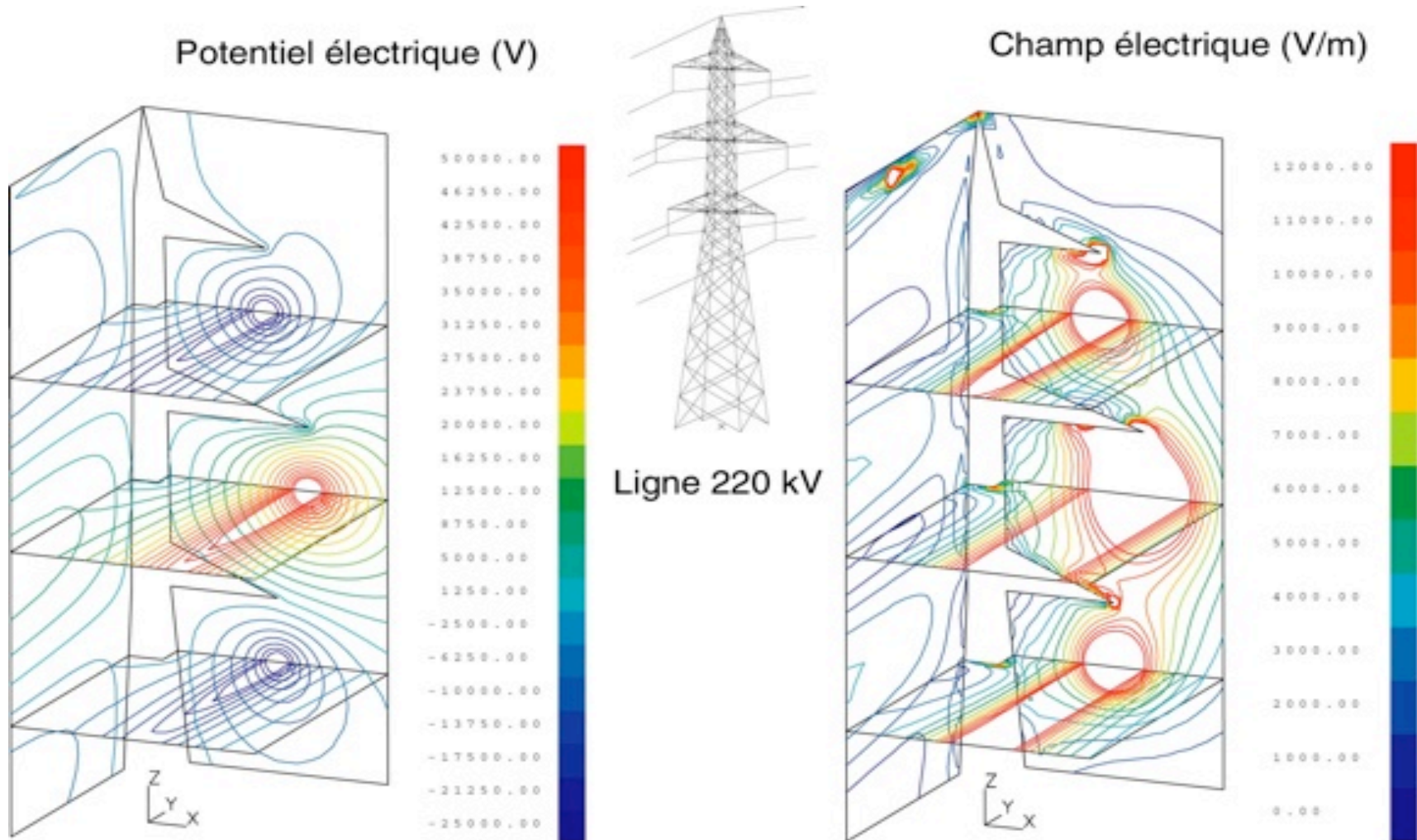


Cable bundles and high-voltage isolators

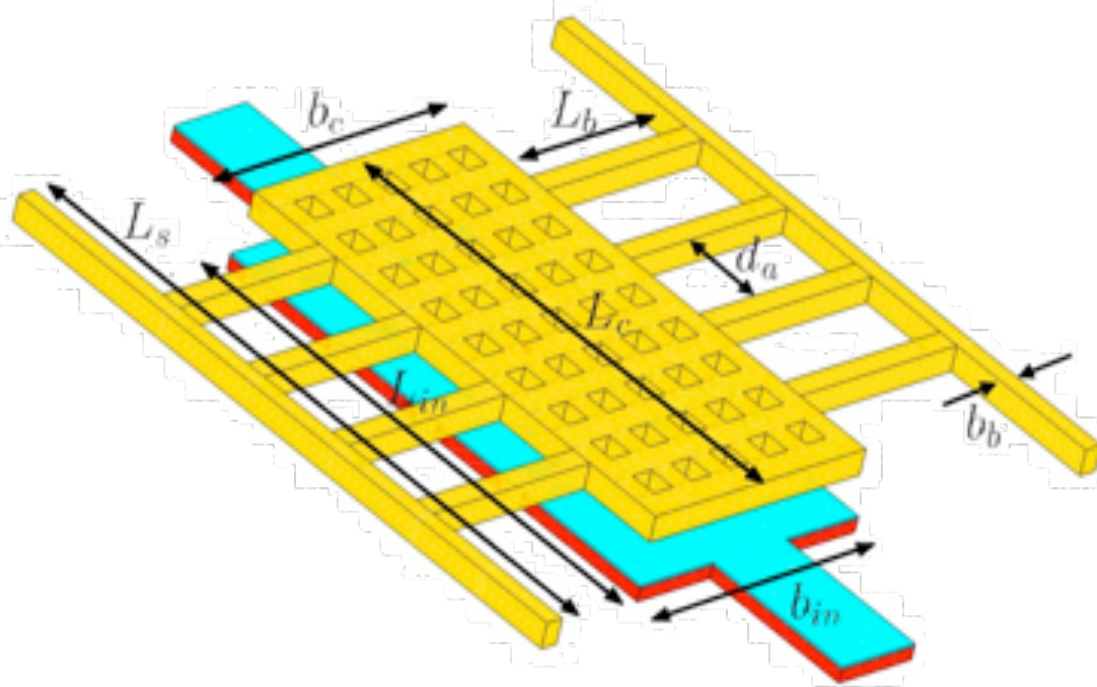


Electrostatics

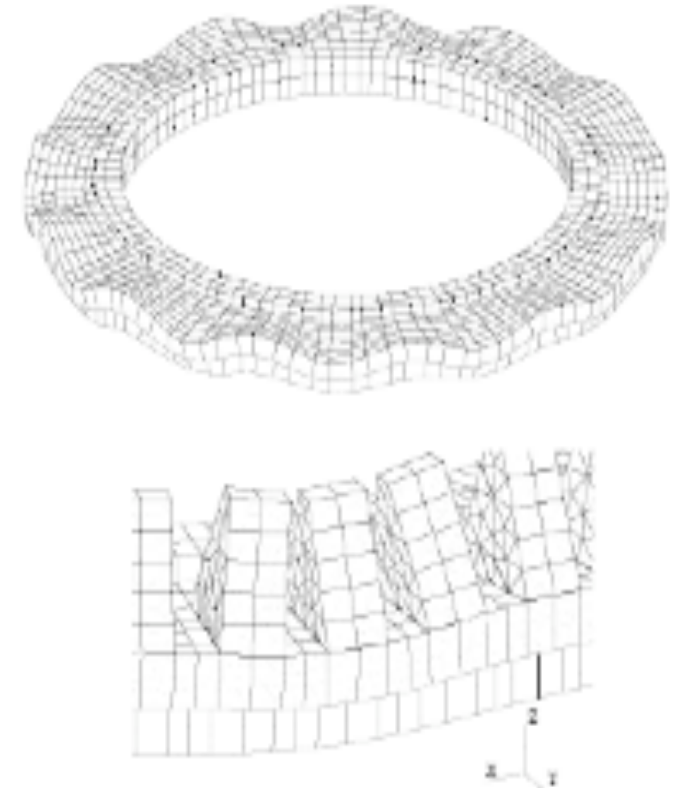
Potential and electric field next to a 220 kV high voltage tower



Electrostatics

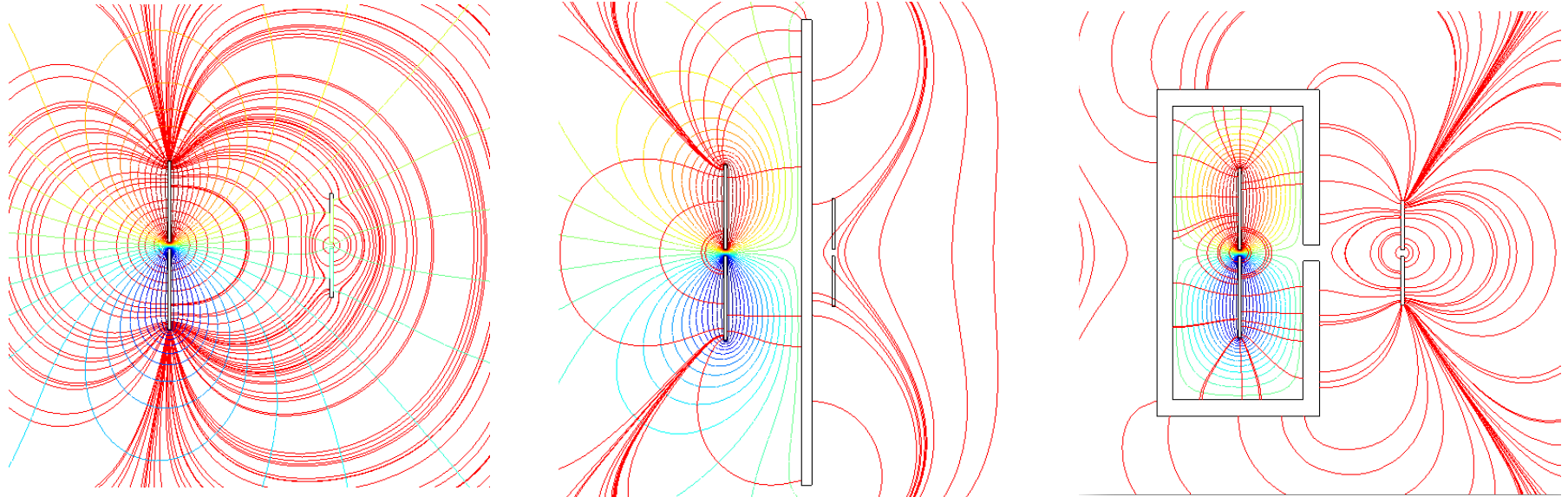


Shunt Capacitive MEM switch



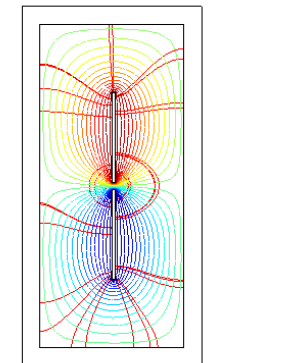
Piezoelectric Motor

Electrostatics



Effectiveness of electric field
shields

(T. Hubing, Clemson University)

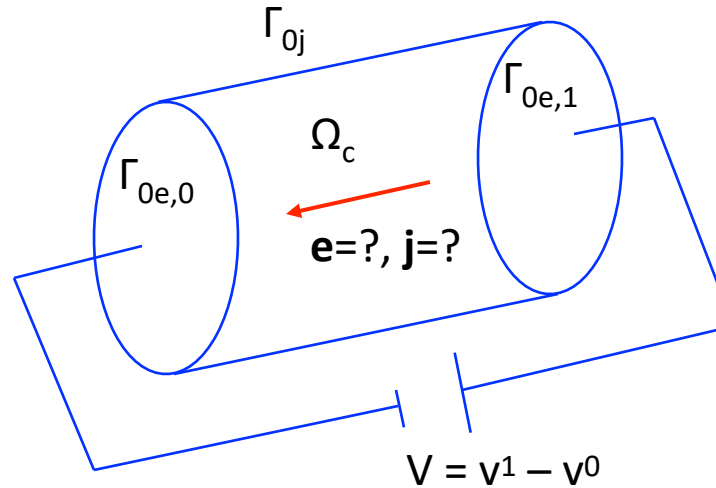


Electrokinetics

$$\mathbf{curl} \mathbf{e} = 0$$

$$\mathbf{curl} \mathbf{h} = \mathbf{j} \Rightarrow \mathbf{div} \mathbf{j} = 0$$

$$\mathbf{j} = \sigma \mathbf{e}$$



Ω_c Conductor

Boundary conditions

$$\mathbf{n} \times \mathbf{e}|_{\Gamma_{0e}} = 0$$

$$\mathbf{n} \cdot \mathbf{j}|_{\Gamma_{0j}} = 0$$

Example: electric scalar potential formulation

$$\mathbf{div} \sigma \mathbf{grad} v = 0 \quad \text{with} \quad \mathbf{e} = -\mathbf{grad} v$$

- Formulation for the conducting region Ω_c
- On each electrode $\Gamma_{0e,i}$, $v = v^i \Rightarrow v|_{\Gamma_{0e,i}} = v^i$

Electro-quasistatics

$$\mathbf{curl} \mathbf{e} = 0$$

$$\mathbf{curl} \mathbf{h} = \mathbf{j} + \partial_t \mathbf{d} \Rightarrow \operatorname{div} (\mathbf{j} + \partial_t \mathbf{d}) = 0$$

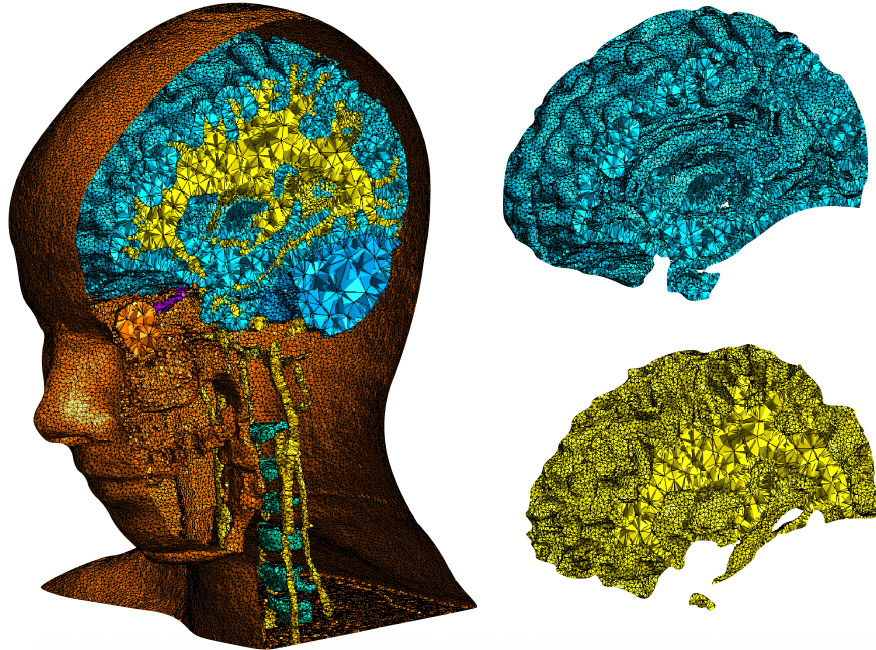
$$\mathbf{j} = \sigma \mathbf{e}$$

$$\mathbf{d} = \varepsilon \mathbf{e}$$

Example: electric scalar potential formulation

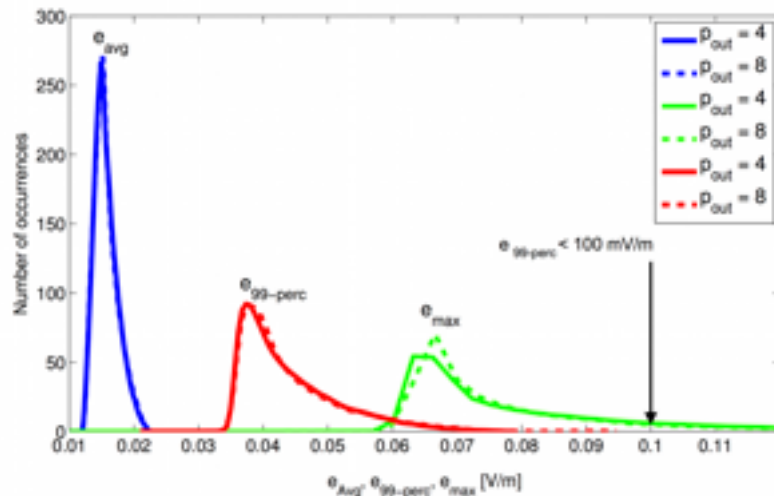
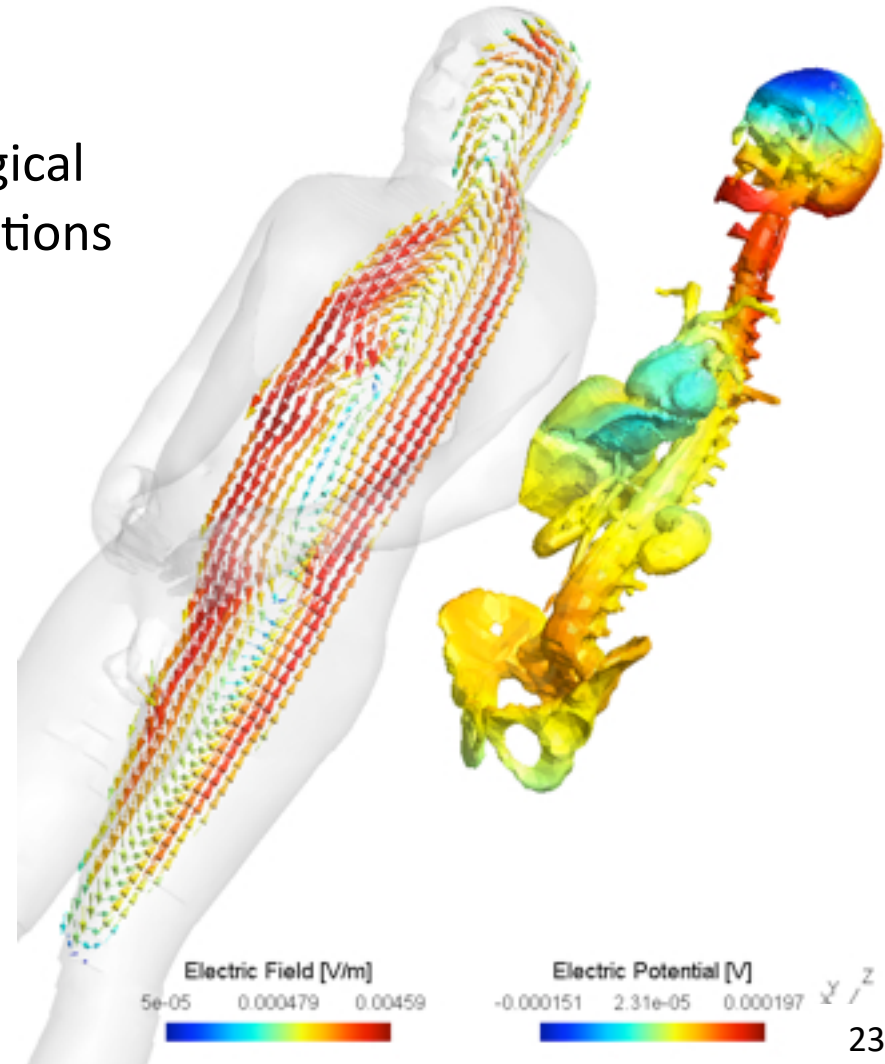
$$\operatorname{div} (\sigma \mathbf{grad} v + \varepsilon \mathbf{grad} \partial_t v) = 0 \quad \text{with} \quad \mathbf{e} = -\mathbf{grad} v$$

Electrokinetics & Electro-quasistatics



Induced Electric Effects due to Magnetic Field from Overhead Power Line
 $B = (100e-6, 0, 0)$ [T]

Biological applications

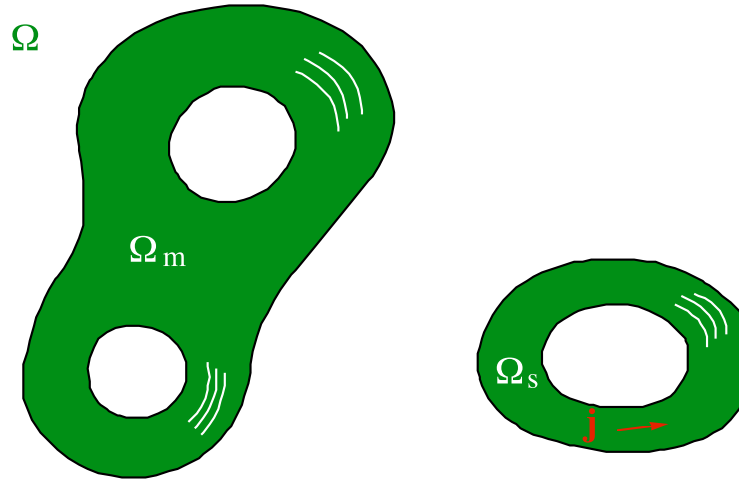


Magnetostatics

$$\text{curl } \mathbf{h} = \mathbf{j}$$

$$\text{div } \mathbf{b} = 0$$

$$\mathbf{b} = \mu \mathbf{h}$$



- Ω Studied domain
- Ω_m Magnetic domain
- Ω_s Inductor

$$\mathbf{j} = \mathbf{j}_s \text{ imposed source current density in inductor}$$

Example: magnetic vector potential formulation

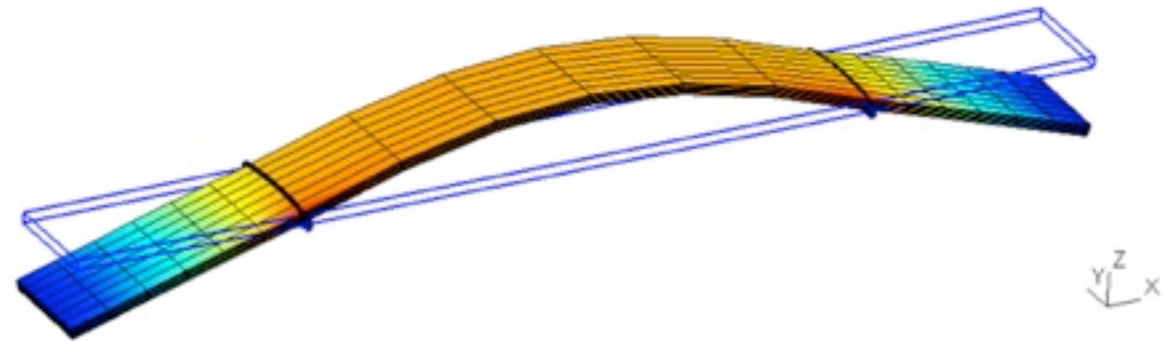
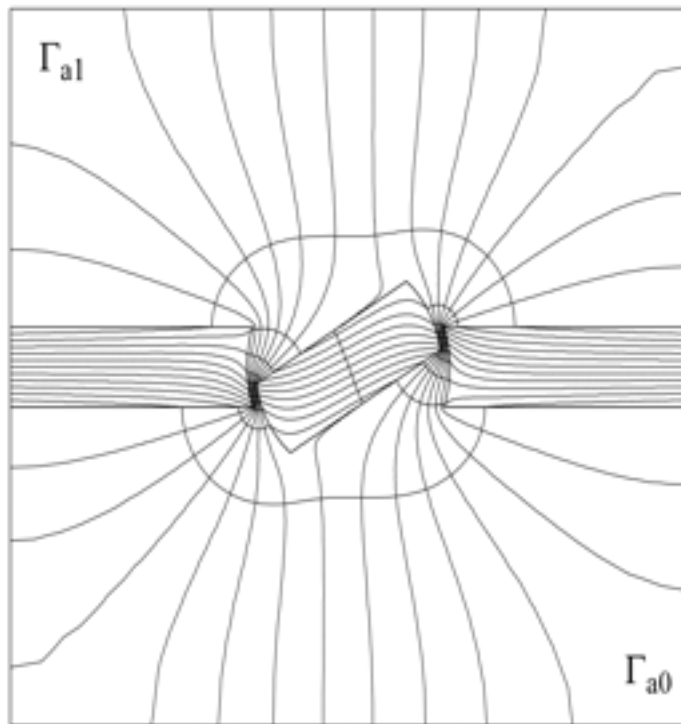
$$\text{curl } \frac{1}{\mu} \text{curl } \mathbf{a} = \mathbf{j}_s \text{ with } \mathbf{b} = \text{curl } \mathbf{a}$$

With magnets:

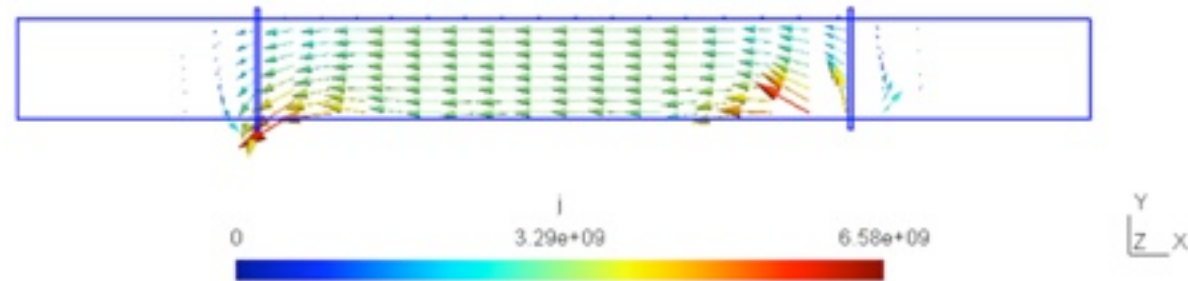
$$\mathbf{b} = \mu \mathbf{h} + \mathbf{b}_s$$

$$\mathbf{h} = \frac{1}{\mu} \mathbf{b} + \mathbf{h}_s$$

Magnetostatics



MEMS magnetometer



Magneto-quasistatics

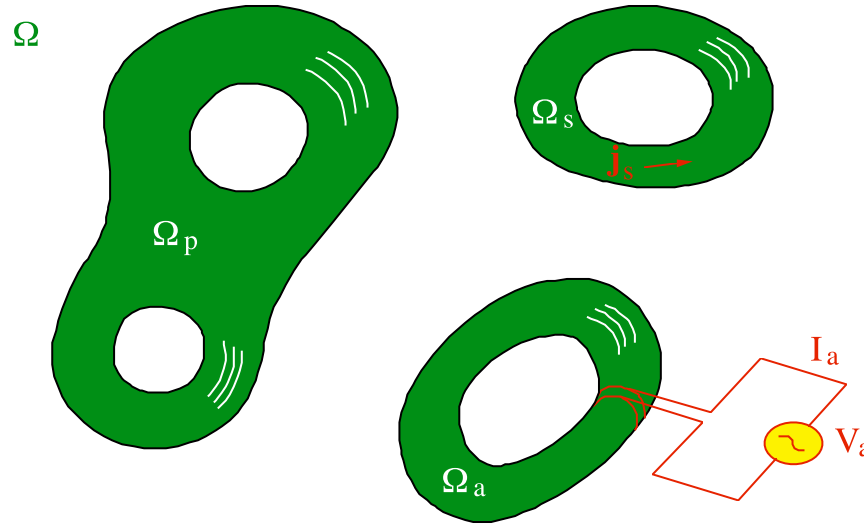
$$\text{curl } \mathbf{h} = \mathbf{j}$$

$$\text{curl } \mathbf{e} = -\partial_t \mathbf{b}$$

$$\text{div } \mathbf{b} = 0$$

$$\mathbf{b} = \mu \mathbf{h} + \mathbf{b}_s$$

$$\mathbf{j} = \sigma \mathbf{e} + \mathbf{j}_s$$



- Ω Studied domain
- Ω_p Passive conductor and/or magnetic domain
- Ω_a Active conductor
- Ω_s Inductor

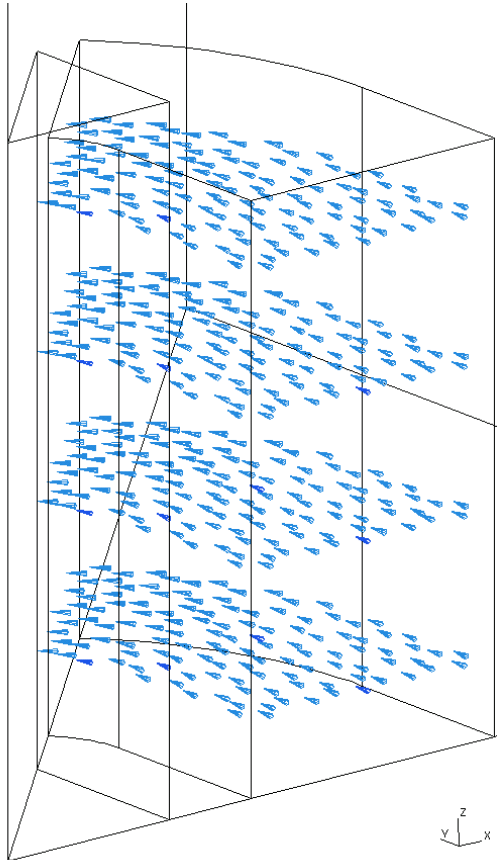
Example: magnetic vector potential formulation

$$\text{curl } \frac{1}{\mu} \text{curl } \mathbf{a} + \sigma (\partial_t \mathbf{a} + \text{grad } v) = \mathbf{j}_s \quad \text{with} \quad \mathbf{b} = \text{curl } \mathbf{a}$$

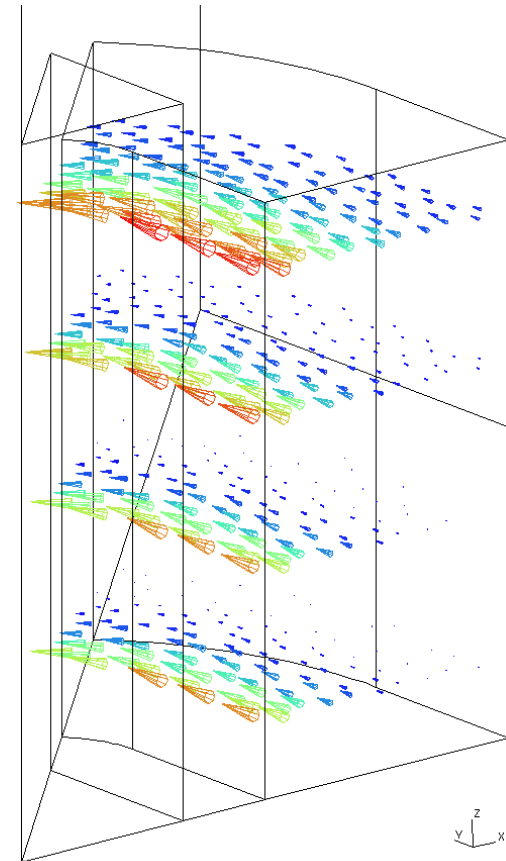
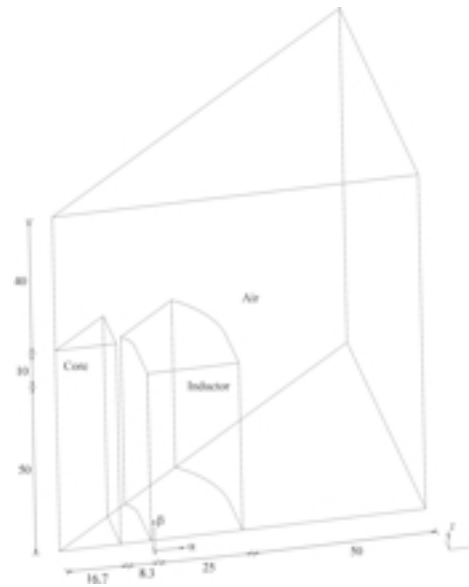
$$\mathbf{e} = -\text{grad } v - \partial_t \mathbf{a}$$

Magneto-quasistatics

Inductor (portion : 1/8th)



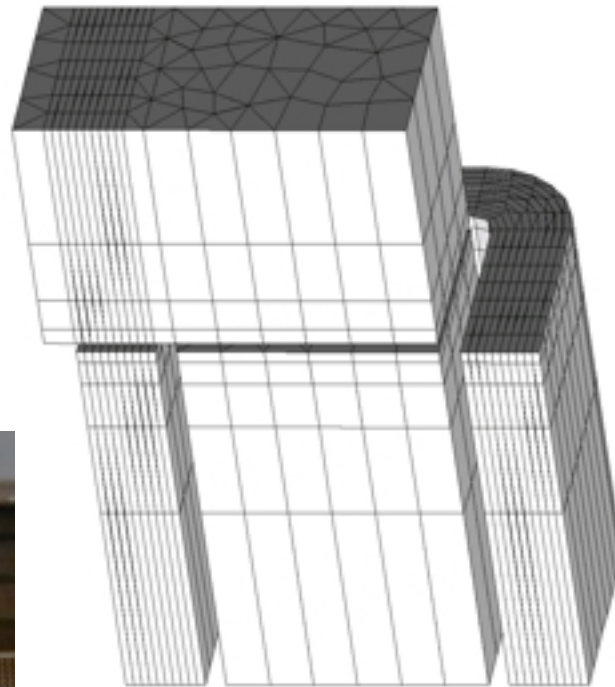
Stranded inductor -
uniform current density (\mathbf{j}_s)



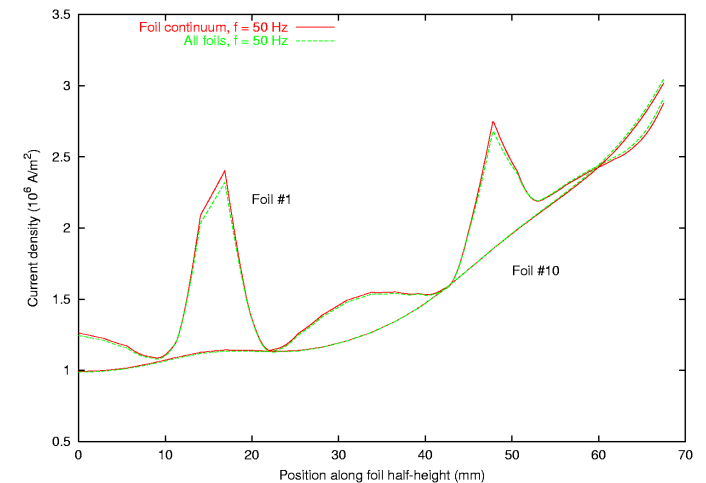
Massive inductor -
non-uniform current density (\mathbf{j})

Magneto-quasistatics

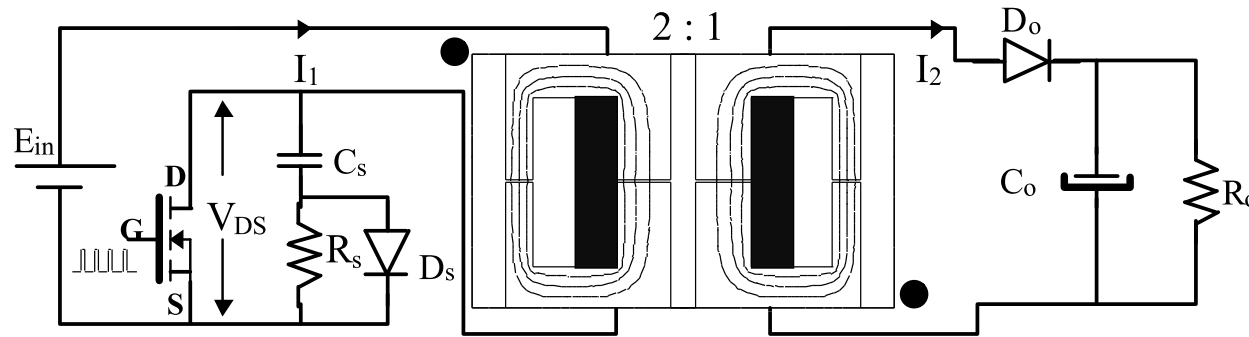
Foil winding inductance - current density (in a cross-section)



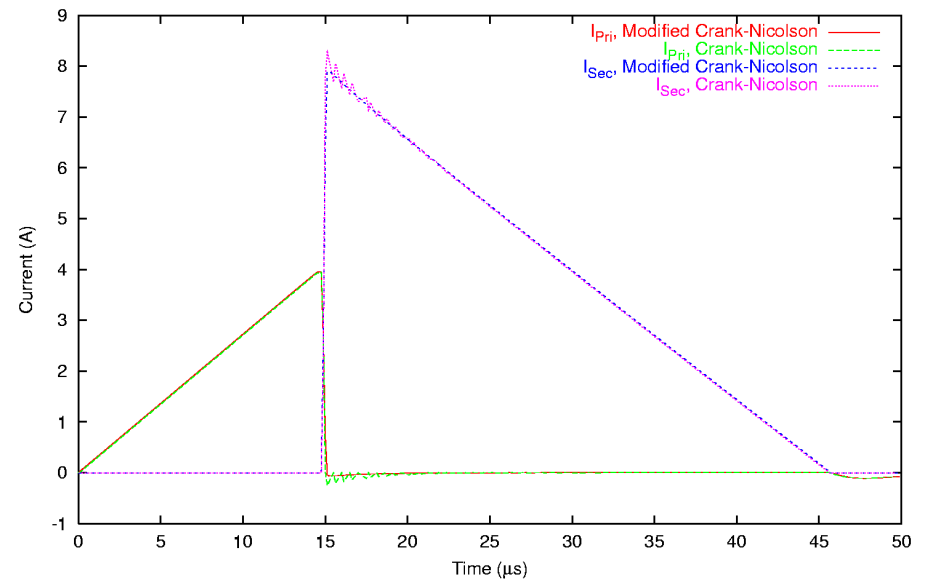
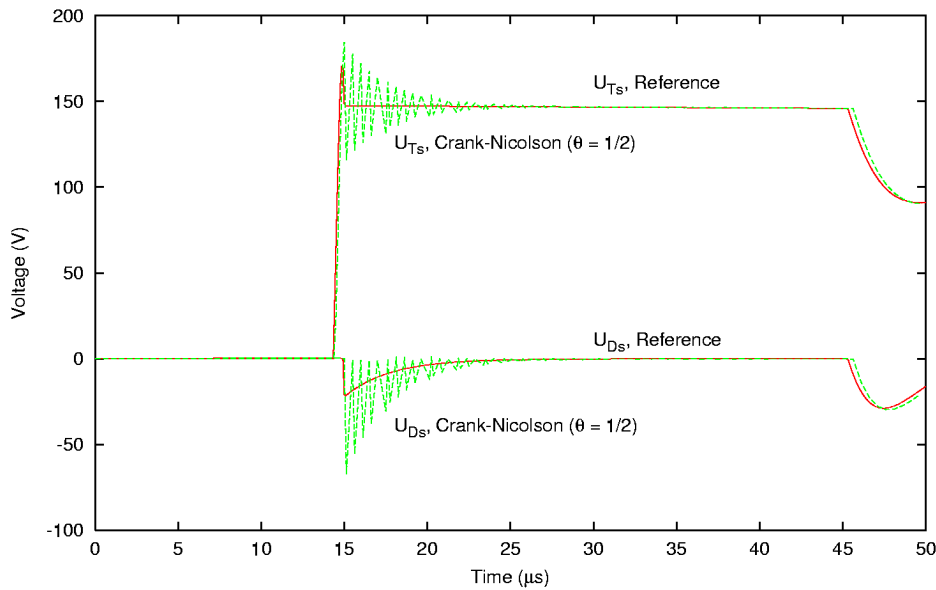
With air gaps, Frequency $f = 50$ Hz



Magneto-quasistatics

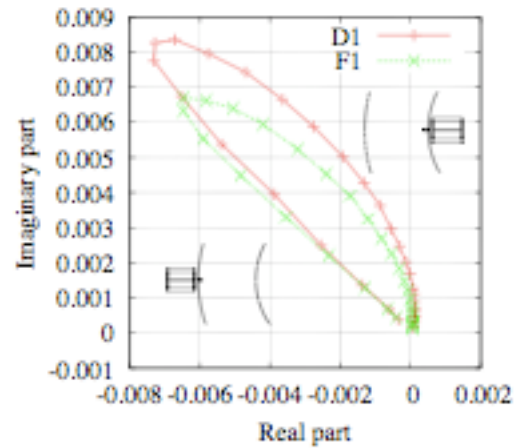
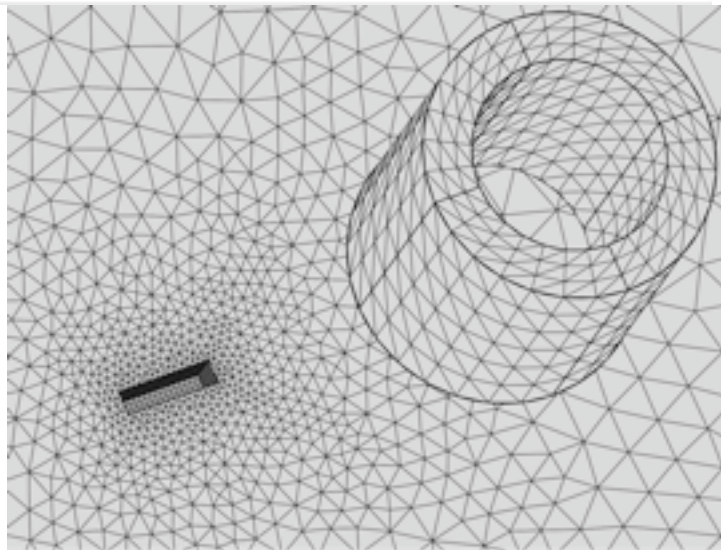
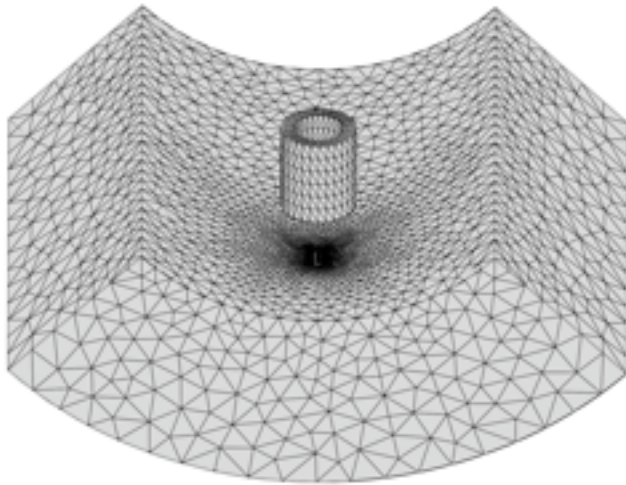


Sudden primary and secondary current changes in the transformer

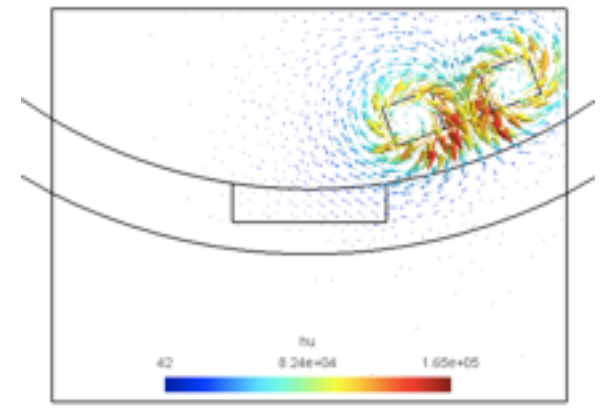


Magneto-quasistatics

Eddy-current non-destructive testing

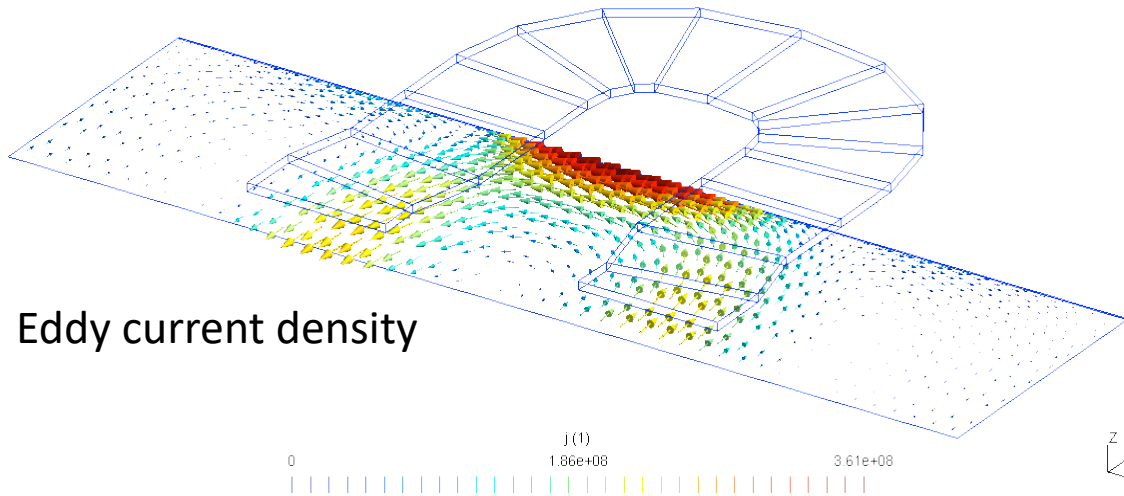


Impedance change

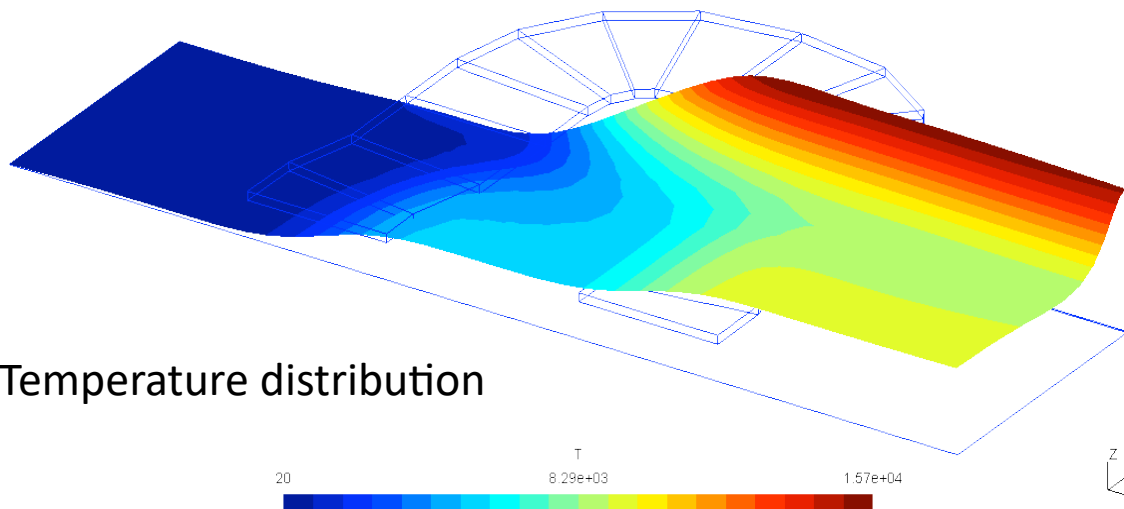


Magnetic field without defect

Magneto-quasistatics



Eddy current density

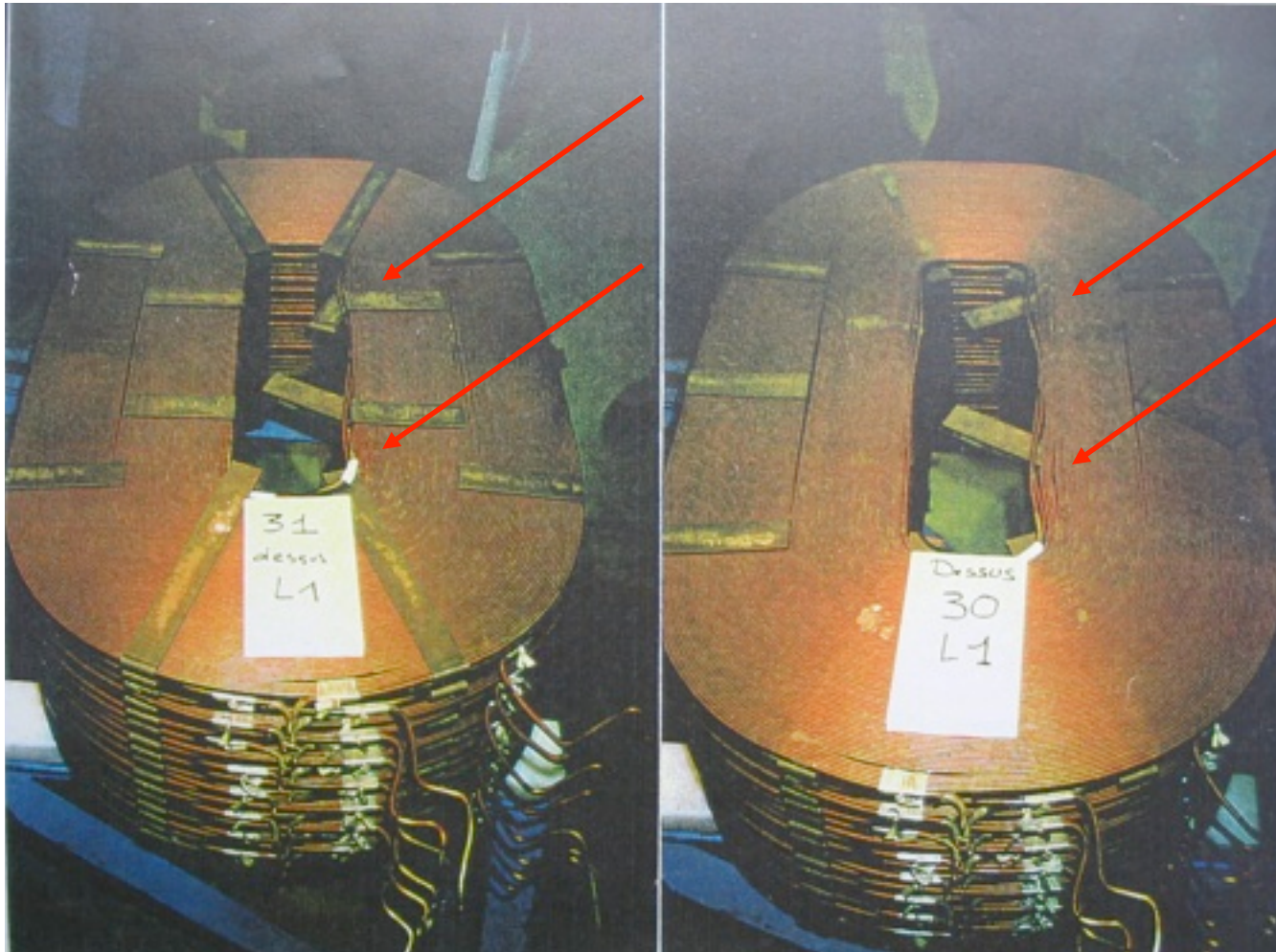


Temperature distribution

Transverse induction heating
(nonlinear physical characteristics,
moving plate, global quantities)

Search for optimization of
temperature profile

Magneto-quasistatics

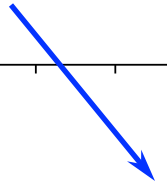
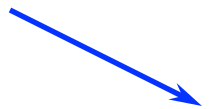
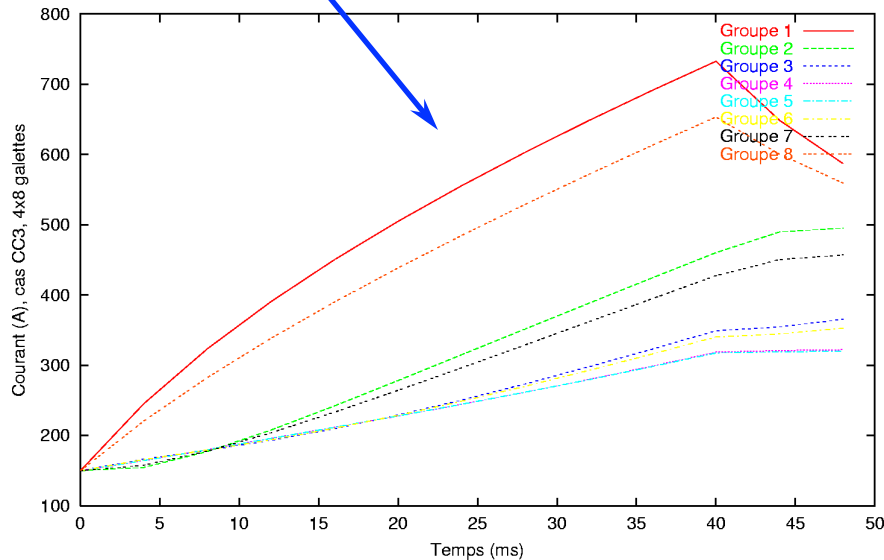


Forces

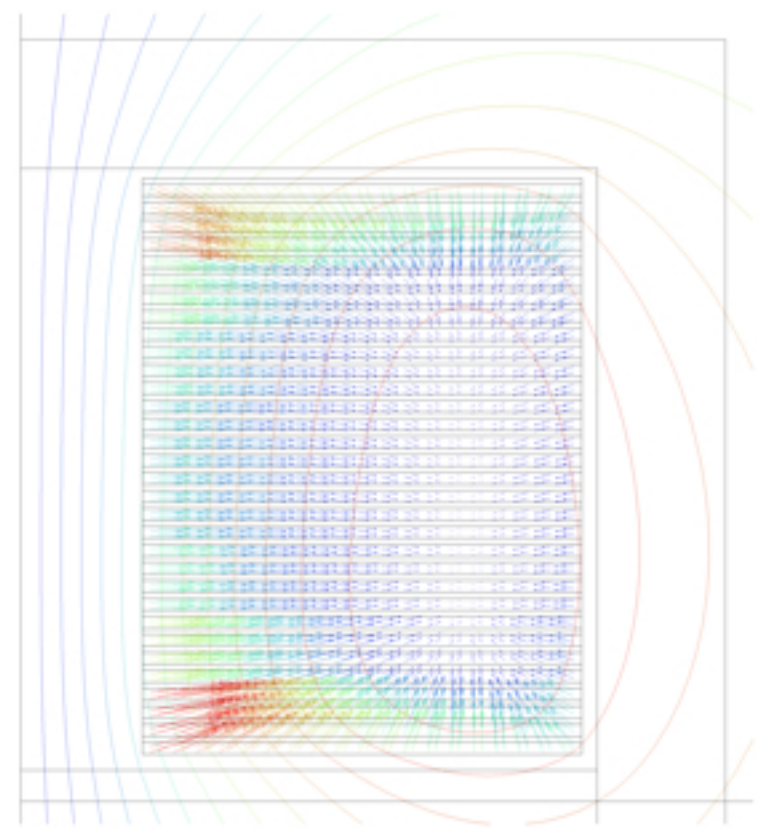
Magneto-quasistatics

Magnetic field lines and electromagnetic force (N/m)
(8 groups, total current 3200 A)

Currents in each of the 8 groups in parallel non-uniformly distributed!



- 1.850e+03
 - 1.718e+03
 - 1.587e+03
 - 1.455e+03
 - 1.324e+03
 - 1.192e+03
 - 1.061e+03
 - 9.293e+02
 - 7.977e+02
 - 6.662e+02
 - 5.347e+02
 - 4.032e+02
 - 2.716e+02
 - 1.401e+02
 - 8.602e+00
- jxb_S



Full Wave

$$\mathbf{curl} \mathbf{h} = \mathbf{j} + \partial_t \mathbf{d}$$

$$\mathbf{curl} \mathbf{e} = -\partial_t \mathbf{b}$$

$$\mathbf{b} = \mu \mathbf{h}$$

$$\mathbf{d} = \varepsilon \mathbf{e}$$

$$\mathbf{j} = \sigma \mathbf{e}$$

+ Silver-Müller radiation condition at infinity (outgoing waves)

Example: electric or magnetic field formulations

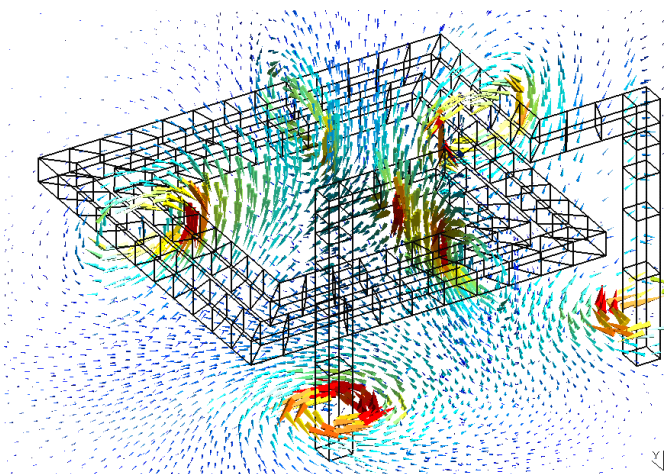
$$\mathbf{curl} \mathbf{curl} \mathbf{e} + \sigma \mu \partial_t \mathbf{e} + \varepsilon \mu \partial_t^2 \mathbf{e} = 0$$

$$\mathbf{curl} \mathbf{curl} \mathbf{h} + \sigma \mu \partial_t \mathbf{h} + \varepsilon \mu \partial_t^2 \mathbf{h} = 0$$

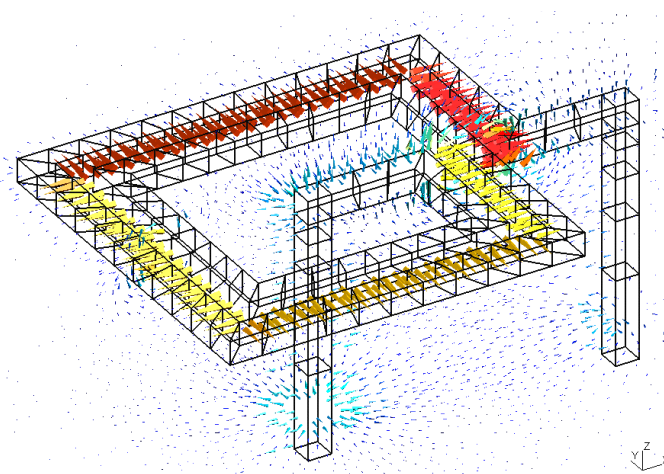
Full Wave

- Frequency and time domain analyses
- Uncoupled resolution

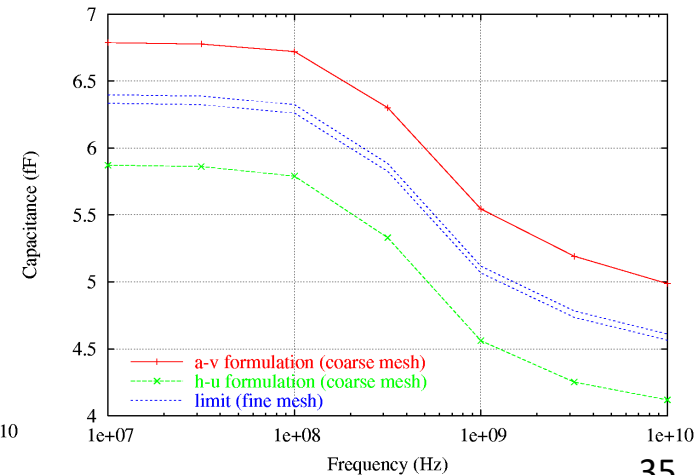
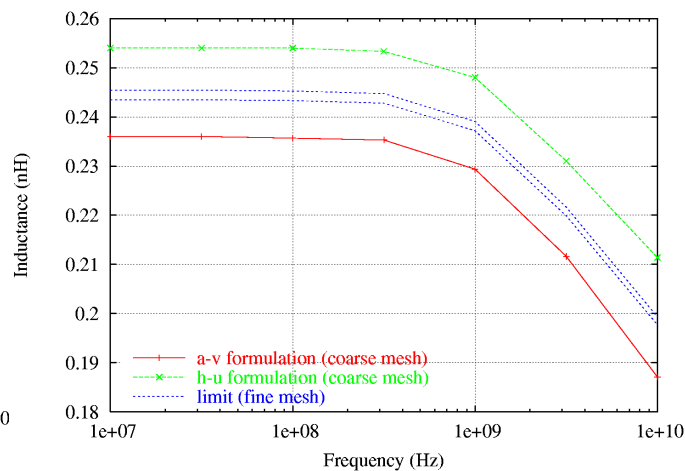
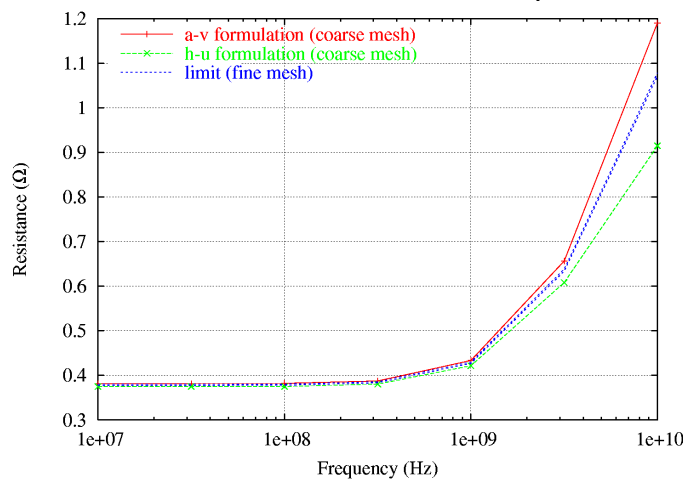
Magnetic flux density



Electric field

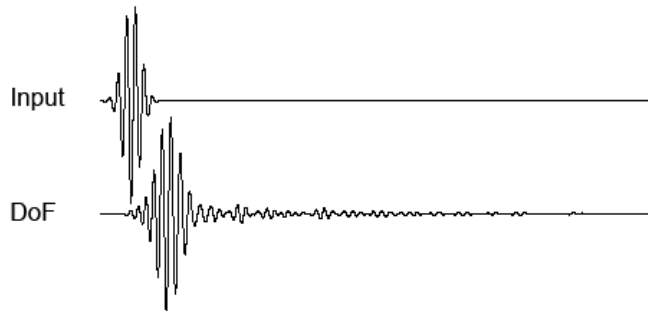
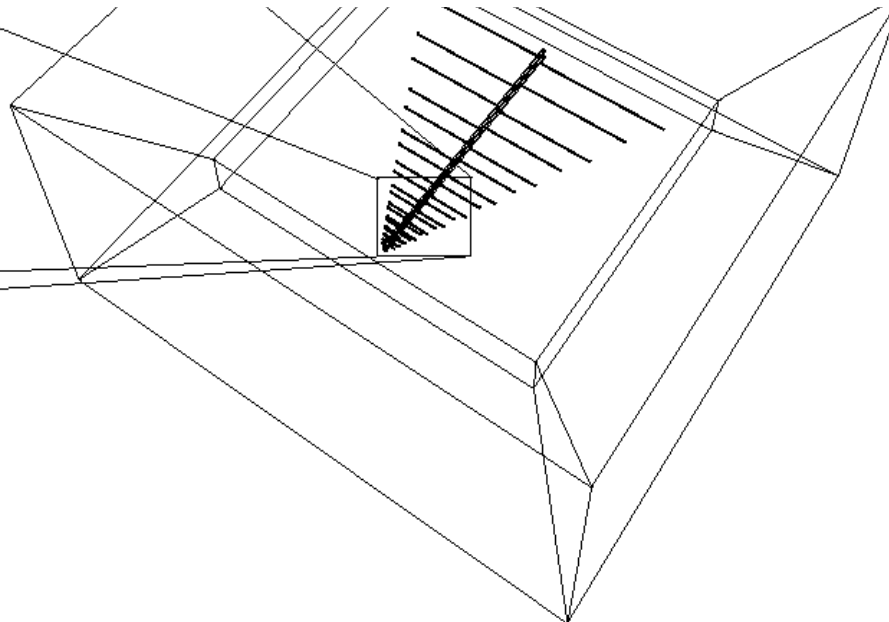
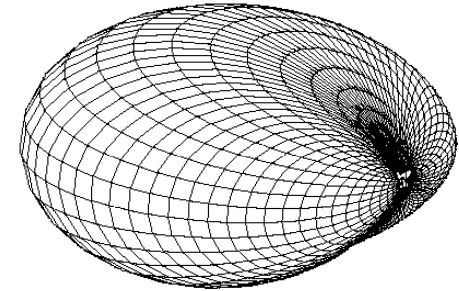
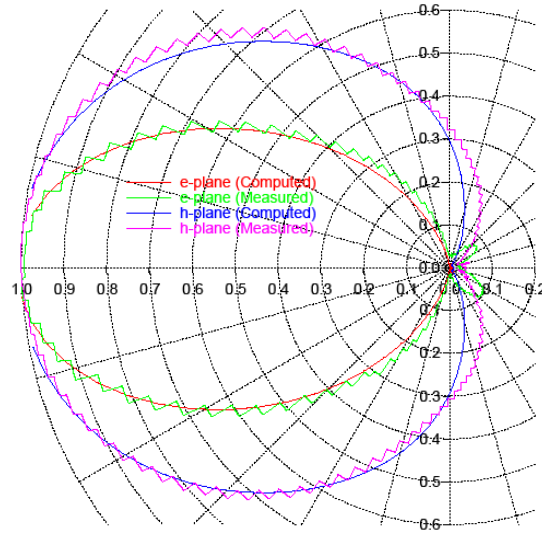
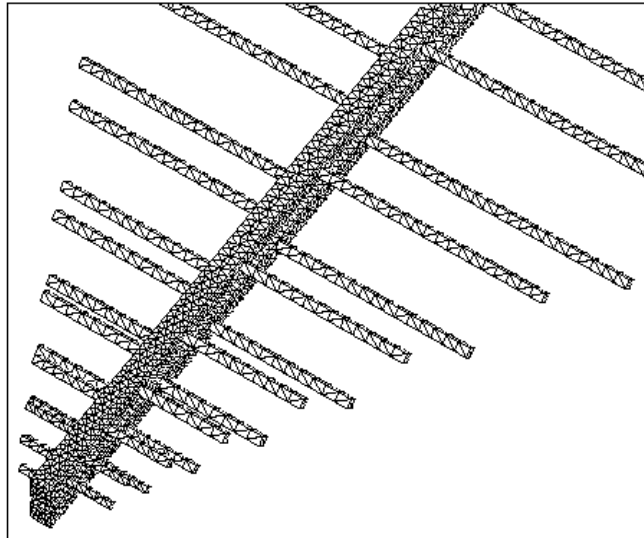


Resistance, inductance and capacitance versus frequency



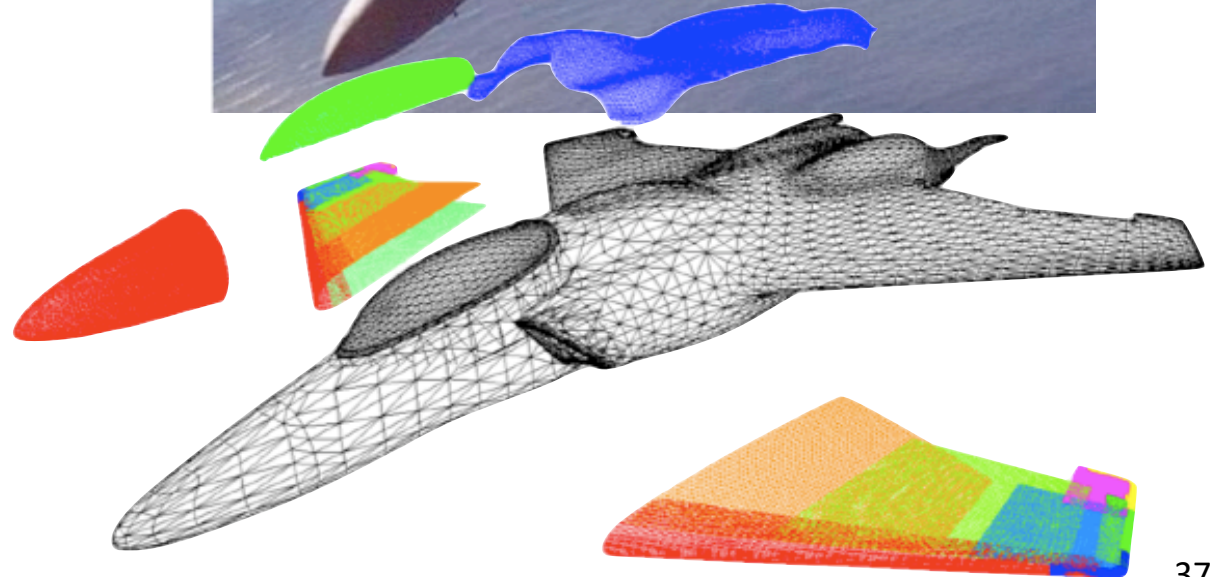
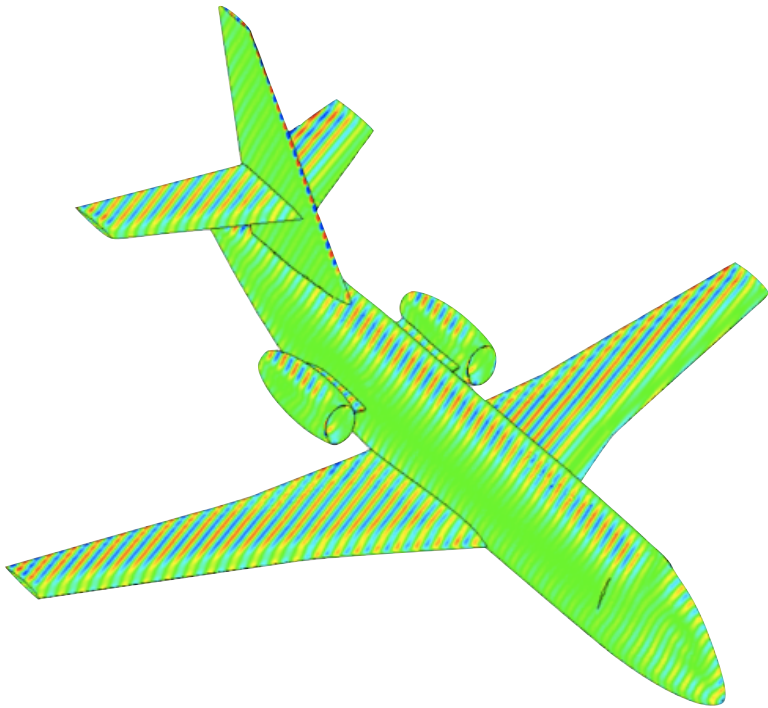
Full Wave

Log-periodic antenna

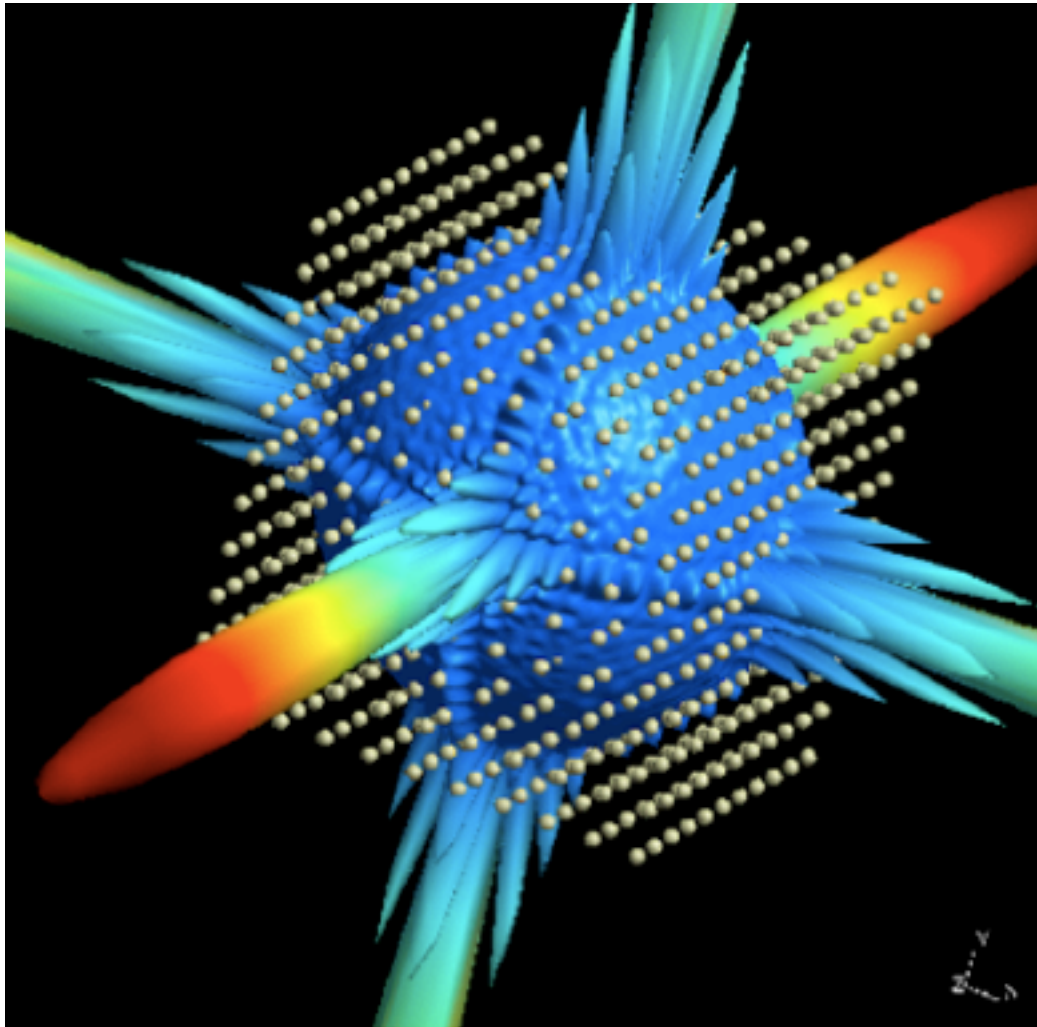


Full Wave

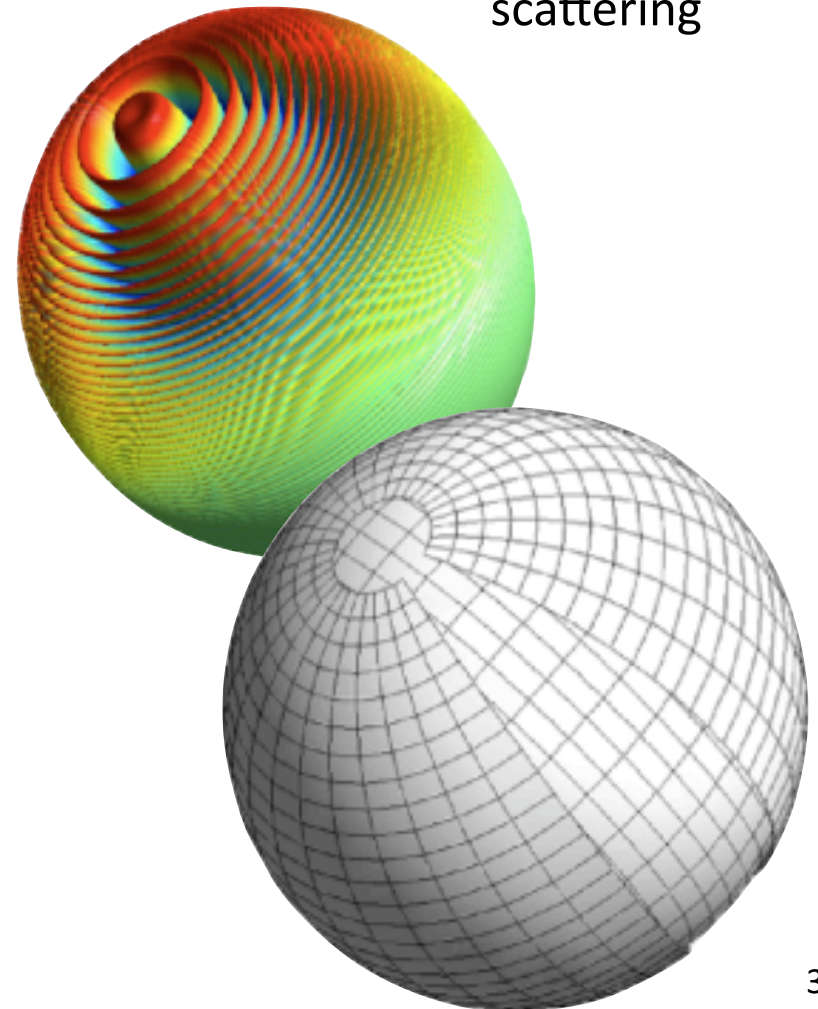
Radar



Full Wave

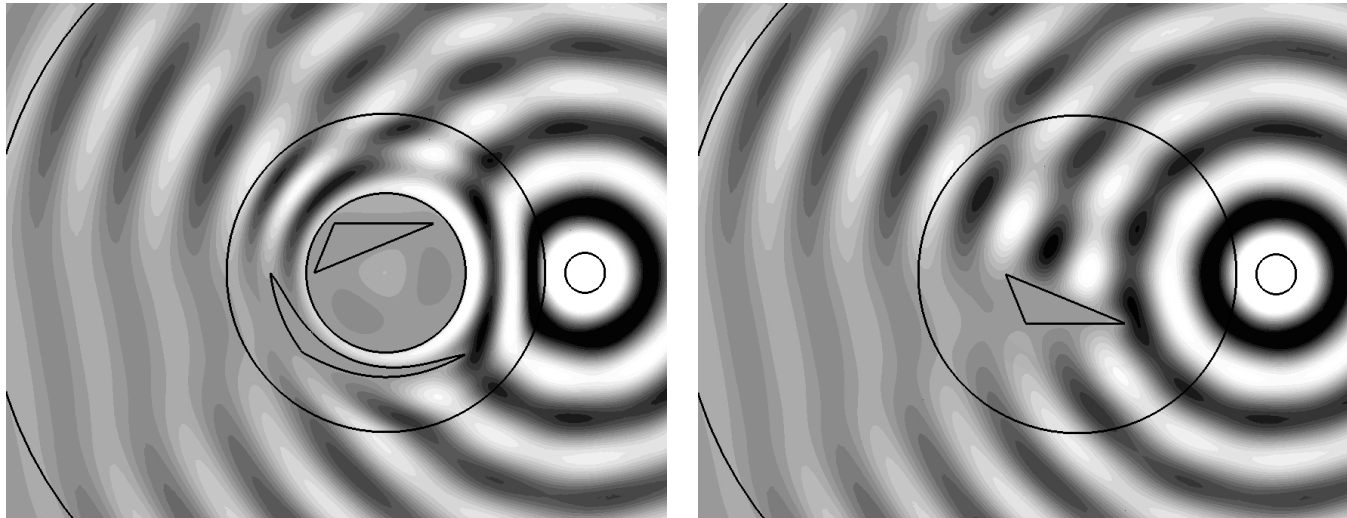


High-frequency
scattering

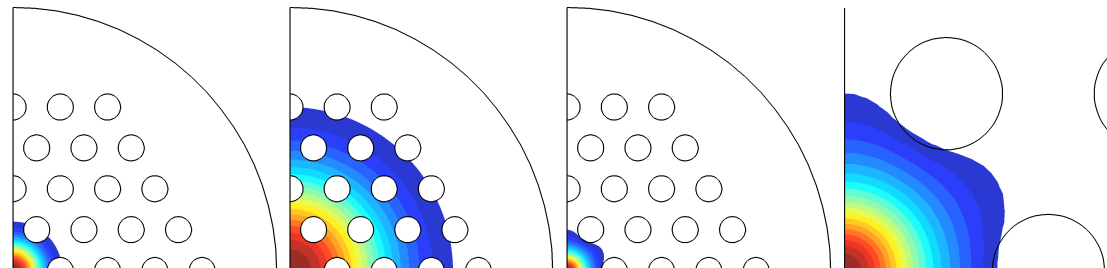
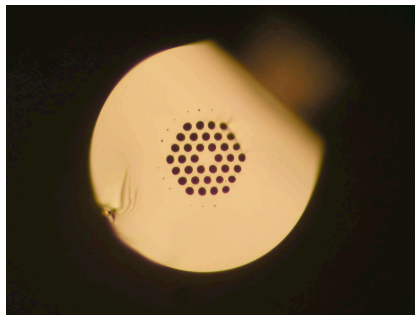


Full Wave

Generalized optical cloaking (“polyjuice”)



Microstructured optical fibers: photonic crystal & non-linear (Kerr) effects



Full Wave

Optical Coherence Tomography (OCT) of human retina

