



#### Lecture 3b Computation of Solutions

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# Numerical solutions of PDEs

- We have found formulae for solutions to several first and second order PDEs, but other problems encountered in practice are not as simple and cannot be solved by formulae
- Even when there is a formula, it might be so complicated that we would prefer to visualize a typical solution by looking at its graph
- Let us start investigating if we can reduce the process of solving a PDE with its auxiliary conditions to a finite number of arithmetic calculations that can be carried out by computer



# Numerical solutions of PDEs

There are dangers in doing so:

- If the numerical method is not carefully designed, the computed solution may substantially differ from the actual one.
- For difficult problems the computation could easily take so long (years...) that it would not be tractable

Two well known numerical methods are Finite Difference (FD) and Finite Element (FE): we will introduce both.

Finite Volume (FV) is another numerical method, widely used in fluid flow computations.



# Learning objectives of this class

Review elementary finite difference (FD) schemes

Become aware the dangers of "blindly" applying FD

Outline

**Review of Finite Difference Schemes** 



Application to the Diffusion Equation





#### 1 – Review of Finite Difference Schemes

### Finite differences

- Finite difference schemes consist in replacing each derivative by a difference quotient.
- Consider a function u(x) of one variable.

Choose a *mesh size*  $\Delta x$ .

Approximate the value  $u(j\Delta x)$  for  $x = j\Delta x$ by a number  $u_j$  indexed by an integer j:

$$u_j \sim u(j\Delta x)$$



## Approximations for the first derivative

The three standard approximations for the first derivative  $\frac{\partial u}{\partial x}(j\Delta x)$  are

The backward difference: 
$$\frac{u_j - u_{j-1}}{\Delta x}$$

The centered difference: 
$$\frac{u_{j+1} - u_{j-1}}{2\Delta x}$$

The forward difference: 
$$\frac{u_{j+1} - u_j}{\Delta x}$$



### Approximations for the first derivative

Each of them is a correct approximation, as shown by a Taylor expansion (assuming *u* is a *C*<sup>4</sup> function):  $u(x + \Delta x) = u(x) + u'(x)\Delta x + \frac{1}{2}u''(x)(\Delta x)^2 + \frac{1}{6}u'''(x)(\Delta x)^3 + O(\Delta x)^4$ 

Replacing  $\Delta x$  by  $-\Delta x$ , we get  $u(x - \Delta x) = u(x) - u'(x)\Delta x + \frac{1}{2}u''(x)(\Delta x)^2 - \frac{1}{6}u'''(x)(\Delta x)^3 + O(\Delta x)^4$ 

We deduce that

$$u'(x) = \frac{u(x) - u(x - \Delta x)}{\Delta x} + O(\Delta x)$$
  

$$\downarrow \text{ Taking } x = j \Delta x$$
  
The backward difference: 
$$\frac{u_j - u_{j-1}}{\Delta x}$$



## Approximations for the first derivative

Similarly for all three standard approximations:

$$u'(x) = \frac{u(x) - u(x - \Delta x)}{\Delta x} + O(\Delta x)$$
  
Backward difference  
$$= \frac{u(x + \Delta x) - u(x)}{\Delta x} + O(\Delta x)$$
  
Forward difference  
$$= \frac{u(x + \Delta x) - u(x - \Delta x)}{2\Delta x} + O(\Delta x)^{2}$$
  
Centred difference

Taking  $x = j\Delta x$  we see that backward and forward differences are correct approximations to the order  $O(\Delta x)$ ; and centered differences is a correct approximation to the order  $O(\Delta x)^2$ 



# Approximations for the second derivative

The simplest approximation for the second derivative is the

centered second difference:  $u''(j\Delta x) \sim \frac{u_{j+1} - 2u_j + u_{j-1}}{(\Delta x)^2}$ 

which is justified by the same Taylor expansions as used before, which when added give:

$$u''(x) = \frac{u(x + \Delta x) - 2u(x) + u(x - \Delta x)}{(\Delta x)^2} + O(\Delta x)^2$$

Centered second difference is thus valid with an error of  $O(\Delta x)^2$ 



Functions of two variables

For a function of 2 variables u(x,t), we choose a mesh size for both variables:

 $u(j\Delta x, n\Delta t) \sim u_j^n$ 

The forward difference approximations of the first order partial derivatives are then for example

$$\frac{\partial u}{\partial t}(j\Delta x, n\Delta t) \sim \frac{u_j^{n+1} - u_j^n}{\Delta t} \qquad \qquad \frac{\partial u}{\partial x}(j\Delta x, n\Delta t) \sim \frac{u_{j+1}^n - u_j^n}{\Delta x}$$



Two types of errors are generally introduced in computations based on such approximations

- 1 Truncation errors refers to the error introduced in the solutions by the approximations themselves, that is, the  $O(\Delta x)$  terms
  - local truncation error: on one term
  - global truncation error: on the actual solution, combining all local contributions
- 2 Round off errors occurs in a real computation because only a certain number of digits, typically 8 or 16, are retained by the computer at each step of the computation







#### 2 – Application to the Diffusion Equation

## Difference equation for Diffusion

Let us solve

$$u_t = u_{xx}, \qquad u(x, 0) = \phi(x)$$

with a forward difference for  $u_t$  and a centered difference for  $u_{xx}$ .

We obtain the following difference equation:

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} = \frac{u_{j+1}^n - 2u_j^n + u_{j-1}^n}{(\Delta x)^2}$$

The local truncation error is  $O(\Delta t)$  for the left-hand side and  $O(\Delta x)^2$  for the right-hand side.



## Difference equation for Diffusion

Let us choose a very small  $\Delta x$  and  $\Delta t = (\Delta x)^2$ .

This leads to the simplified equation:

$$u_{j}^{n+1} = u_{j+1}^{n} - u_{j}^{n} + u_{j-1}^{n}$$

Consider the following initial data  $\phi(x)$ 





# Difference equation for Diffusion

Applying the scheme we obtain:

This is *clearly* not correct: the initial data gets amplified and oscillates – which is not what we expect for diffusion.



We will analyze this in detail next week.

### Take-home messages

- Finite difference approximations can be used to find approximate solutions to PDEs
- Small truncation errors do not guarantee that the solution will be close to the true solution

# Next lecture

The diffusion equation:

1. Theory



2. Detailed stability analysis of FD approximation